Introduction to Water Simulation II ---Numerical Solution

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Outline

- Splitting the Navier-Stokes Equation
- MAC Grid, a staggered grid
- Algorithms details
 - Advection
 - Add body forces
 - Make water incompressible
- From velocity field to water surface

A close look at the Navier-Stokes Equation

$$\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} + \frac{1}{\rho} \nabla p = \vec{g} + \nu \nabla \cdot \nabla \vec{u}$$
$$\overset{\parallel}{\frac{D\vec{u}}{Dt}} \nabla \cdot \vec{u} = 0$$

- Unknowns: \vec{u} and p. \vec{u} is what we want.
- $\blacktriangleright p$ is constrained by the incompressibility term
- Viscosity term could be dropped
- Still too complicated to solve it in one step. It would be nice to split it into several steps.

A close look at the Navier-Stokes Equation

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Dived and Conquer

A simple example

$$\frac{dq}{dt} = f(q) + g(q)$$

Euler Solution

 $q^{n+1} = q^n + \Delta t \left[f(q) + g(q) \right]$

Split it into two ODEs

$$\frac{dq}{dt} = f(q)$$
$$\frac{dq}{dt} = g(q)$$

Solve them sequentially

$$\tilde{q} = q^n + \Delta t f(q^n)$$
$$q^{n+1} = \tilde{q} + \Delta t g(\tilde{q})$$

Splitting the Fluid Equations

$$\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} + \frac{1}{\rho} \nabla p = \vec{g}$$

$$\prod_{\substack{\mathbf{i} \\ f_1(\vec{u}) \\ \mathbf{i} \\ f_2(\vec{u}) \\ f_$$

$$\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} = 0 \qquad \Longrightarrow \qquad \frac{D \vec{u}}{D t} = 0$$

$$\qquad \qquad \frac{\partial \vec{u}}{\partial t} = \vec{g}$$

Body Forces

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Pressure/Incompressiblity

s.t. $\nabla \cdot \vec{u} = 0$

Splitting the Fluid Equations

$$\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} + \frac{1}{\rho} \nabla p = \vec{g}$$
$$\prod_{f_1(\vec{u})}^{"} f_3(\vec{u}) = f_2(\vec{u})$$

- Start with an initial divergence-free velocity field $\vec{u}^{(0)}$.
- For time step n = 0, 1, 2, ...
 - Determine a good time step Δt to go from time t_n to time t_{n+1}
 - Set $\vec{u}^A = \text{advect}(\vec{u}^n, \Delta t, \vec{u}^n)$
 - Add $\vec{u}^B = \vec{u}^A + \Delta t \vec{g}$
 - Set $\vec{u}^{n+1} = \text{project}(\Delta t, \vec{u}^B)$

 $\begin{aligned} \frac{D\vec{u}}{Dt} &= 0\\ \frac{\partial \vec{u}}{\partial t} &= \vec{g}\\ \frac{\partial \vec{u}}{\partial t} + \frac{1}{\rho} \nabla p &= 0 \quad \text{ s.t. } \quad \nabla \cdot \vec{u} = 0 \end{aligned}$

Discretization In Space

MAC Grid (staggered grid)



- Benefits
 - Accurate central difference
- Downsides
 - Variables spread up, interpolation needed.

The Disaster of Simple Grid When Evaluating the Derivatives

forward or backward difference

$$\left(\frac{\partial q}{\partial x}\right)_i \approx \frac{q_{i+1} - q_i}{\Delta x}$$

central difference

$$\left(\frac{\partial q}{\partial x}\right)_i \approx \frac{q_{i+1} - q_{i-1}}{2\Delta x}$$

A bad situation of central difference, zero derivative everywhere.



The Disaster of Simple Grid When Evaluating the Derivatives

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central difference

$$\left(\frac{\partial q}{\partial x}\right)_i \approx \frac{q_{i+1} - q_{i-1}}{2\Delta x}$$



3D MAC Grid



- For each Cell
 - pressure locates at the center
 - Each component of the velocity takes up two faces
 - Each face only has one component of the velocity, interpolation needed.
- Derivatives of velocity at the center of the cell...
- Derivatives of pressure at the center of each facet...

3D MAC Grid



Interpolate the velocity at the center of the cell and its facets.

$$\vec{u}_{i,j,k} = \left(\frac{u_{i-1/2,j,k} + u_{i+1/2,j,k}}{2}, \frac{v_{i,j-1/2,k} + v_{i,j+1/2,k}}{2}, \frac{w_{i,j,k-1/2} + w_{i,j,k+1/2}}{2}\right)$$

$$\vec{u}_{i+1/2,j,k} = \left(\begin{array}{c} v_{i,j-1/2,k} + v_{i,j+1/2,k} & w_{i,j,k-1/2} + w_{i,j,k+1/2} \\ + v_{i+1,j-1/2,k} + v_{i+1,j+1/2,k} & \frac{w_{i,j,k-1/2} + w_{i,j,k+1/2}}{4} \\ \end{array}\right)$$

$$\vec{u}_{i,j+1/2,k} = \left(\begin{array}{c} u_{i-1/2,j,k} + u_{i+1/2,j,k} & w_{i,j+1/2,k} \\ + u_{i-1/2,j+1,k} + u_{i+1/2,j+1,k} & v_{i,j+1/2,k} \\ \frac{u_{i-1/2,j,k} + u_{i+1/2,j+1,k}}{4}, & v_{i,j+1/2,k} & \frac{w_{i,j-1/2,k} + v_{i,j+1/2,k}}{4} \\ \end{array}\right)$$

$$\vec{u}_{i,j,k+1/2} = \left(\begin{array}{c} u_{i-1/2,j,k} + u_{i+1/2,j,k} & v_{i,j-1/2,k} + v_{i,j+1/2,k} \\ \frac{u_{i-1/2,j,k} + u_{i+1/2,j,k+1}}{4}, & \frac{v_{i,j-1/2,k+1} + v_{i,j+1/2,k+1}}{4}, & w_{i,j,k+1/2} \\ \end{array}\right)$$

STEP I: Advect the Velocity

$$\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} = 0 \quad \text{or} \quad \frac{D\vec{u}}{Dt} = 0$$

$$\frac{\partial \vec{u}}{\partial t} = \vec{g}$$

$$\frac{\partial \vec{u}}{\partial t} + \frac{1}{\rho} \nabla p = 0 \quad \text{s.t.} \quad \nabla \cdot \vec{u} = 0$$

Advecting Quantities

The goal is to solve

$$\frac{Dq}{Dt} = 0$$

"the advection equation" for any grid quantity q

- advect each component of velocity separately
- Intead of treating it in Euler fashion by directly solving

$$\frac{\partial q}{\partial t} + u \frac{\partial q}{\partial x} = 0$$

we are using Lagrangian notion.

We're on an Eulerian grid, though, so the result will be called "semi-Lagrangian".

Proposed by Jos Stam, "Stable Fluids", 1999

Semi-Lagrangian Algorithm

- For each grid point, find X^{old}, and use the quantity (of previous time step) at this position as the new quantity of the grid point.
- Interpolation may be necessary. Be careful when doing interpolation in staggered grid.
- Forward Euler is adequate to find the old position.

$$\vec{x}_q = \vec{x}_p - \Delta t \vec{u}_p$$



Boundary Conditions

- What if the particle flies out of the water boundary
 - due to numeric error, just extrapolate from nearest points on the boundary;
 - due to water flowing in from outside, ...



Dissipation

- Interpolation cause smoothed velocity field. Small vortices will be phased out.
- It equals to simulate a fluid with viscosity.
- Will be covered by the following lecture.

STEP II: Add Body Forces

$$\begin{aligned} & \frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} = 0 \quad \text{or} \quad \frac{D \vec{u}}{D t} = 0 \\ & \frac{\partial \vec{u}}{\partial t} = \vec{g} \end{aligned}$$

$$\begin{aligned} & \frac{\partial \vec{u}}{\partial t} + \frac{1}{\rho} \nabla p = 0 \quad \text{s.t.} \quad \nabla \cdot \vec{u} = 0 \end{aligned}$$

Integrating Body Forces

- Supper Easy!!!
- Just add the new term at each grid point

$$\vec{u}^* = \vec{u}^{advected} + \Delta t \vec{g}$$

STEP III: Making Fluid Incompressible

$$\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} = 0 \quad \text{or} \quad \frac{D\vec{u}}{Dt} = 0$$

$$\frac{\partial \vec{u}}{\partial t} = \vec{g}$$

$$\frac{\partial \vec{u}}{\partial t} + \frac{1}{\rho} \nabla p = 0 \quad \text{s.t.} \quad \nabla \cdot \vec{u} = 0$$

D

 Space is continuous, but still assume the time space is discrete. Update the velocity,

$$\vec{u}^{n+1} = \vec{u}^* - \frac{\Delta t}{\rho} \nabla p$$

> To make it incompressible, the divergence should be zero.

$$\nabla \cdot (\vec{u}^* - \frac{\Delta t}{\rho} \nabla p) = 0$$
$$-\frac{\Delta t}{\rho} \nabla \cdot \nabla p = -\nabla \cdot \vec{u}^* \qquad \text{Laplace Equation}$$

 $\frac{\partial \vec{u}}{\partial t} + \frac{1}{\rho} \nabla p = 0$

s.t. $\nabla \cdot \vec{u} = 0$

 Solve Laplace Equation to get p, then substitute p into update equation. (Right now, let's leave out boundary conditions, and assume the water is boundless.)

The Discrete Version

- $\frac{\partial \vec{u}}{\partial t} + \frac{1}{\rho} \nabla p = 0$
s.t. $\nabla \cdot \vec{u} = 0$
- For clarity, discretize the pressure equation and the divergence constraint instead.

$$\begin{aligned} \text{discretize } \vec{u}^{n+1} &= \vec{u}^* - \frac{\Delta t}{\rho} \nabla p \\ \text{discretize } \nabla \cdot \vec{u} &= \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \\ \\ u_{i+1/2,j,k}^{n+1} &= u_{i+1/2,j,k} - \Delta t \frac{1}{\rho} \frac{p_{i+1,j,k} - p_{i,j,k}}{\Delta x} \\ v_{i,j+1/2,k}^{n+1} &= v_{i,j+1/2,k} - \Delta t \frac{1}{\rho} \frac{p_{i,j+1,k} - p_{i,j,k}}{\Delta x} \\ w_{i,j,k+1/2}^{n+1} &= w_{i,j,k+1/2} - \Delta t \frac{1}{\rho} \frac{p_{i,j,k+1} - p_{i,j,k}}{\Delta x} \\ \end{aligned}$$

$$\begin{aligned} \frac{\partial \vec{u}}{\partial t} + \frac{1}{\rho} \nabla p &= 0 \\ \text{S.t. } \nabla \cdot \vec{u} &= 0 \end{aligned}$$

$$\begin{aligned} & \text{discretize } \vec{u}^{n+1} = \vec{u}^* - \frac{\Delta t}{\rho} \nabla p \\ & \text{discretize } \vec{u}^{n+1} = \vec{u}^* - \frac{\Delta t}{\rho} \nabla p \\ & \text{discretize } \nabla \cdot \vec{u} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \\ \hline u_{i+1/2,j,k}^{n+1} &= u_{i+1/2,j,k} - \Delta t_i \frac{1}{\rho} \frac{p_{i+1,j,k} - p_{i,j,k}}{\Delta x} \\ u_{i,j+1/2,k}^{n+1} &= w_{i,j+1/2,k} - \Delta t_i \frac{1}{\rho} \frac{p_{i,j+1,k} - p_{i,j,k}}{\Delta x} \\ w_{i,j,k+1/2}^{n+1} &= w_{i,j,k+1/2} - \Delta t_i \frac{1}{\rho} \frac{p_{i,j,k+1} - p_{i,j,k}}{\Delta x} \\ w_{i,j,k+1/2}^{n+1} &= w_{i,j,k+1/2} - \Delta t_i \frac{1}{\rho} \frac{p_{i,j,k+1} - p_{i,j,k}}{\Delta x} \\ & \text{Substitute left equations into the right one} \\ \frac{\Delta t}{\rho} \begin{pmatrix} \frac{6p_{i,j,k} - p_{i+1,j,k} - p_{i,j,k+1}}{\Delta x^2} \\ 0 \end{pmatrix} \\ & - \begin{pmatrix} \frac{u_{i+1/2,j,k} - u_{i-1/2,j,k}}{\Delta x} + \frac{v_{i,j+1/2,k} - v_{i,j-1/2,k}}{\Delta x} \\ 0 \end{pmatrix} \\ & - \begin{pmatrix} \frac{\Delta t}{\rho} \nabla \cdot \nabla p = -\nabla \cdot \vec{u}^* \end{pmatrix} \end{aligned}$$

Putting Them In Matrix-Vector Form

• Each cell has such a linear equation, combining them together we could get p at each cell.

$$\frac{\Delta t}{\rho} \left(\frac{\frac{6p_{i,j,k} - p_{i+1,j,k} - p_{i,j+1,k} - p_{i,j,k+1}}{\Delta x^2}}{\Delta x^2} \right) = - \left(\frac{\frac{u_{i+1/2,j,k} - u_{i-1/2,j,k}}{\Delta x} + \frac{v_{i,j+1/2,k} - v_{i,j-1/2,k}}{\Delta x}}{\Delta x} + \frac{\frac{w_{i,j,k+1/2} - w_{i,j,k-1/2}}{\Delta x}}{\Delta x} \right)$$

 Represent p of all the cells as a linear vector. Write all the linear equations into a matrix-vector form

$$\begin{bmatrix} \ddots & & & & & \\ \cdots & -1 & \cdots & -1 & 6 & -1 & \cdots & -1 & \cdots \\ & & & & \ddots \end{bmatrix} \begin{vmatrix} p_{i-I,j,k} \\ \vdots \\ p_{i,j,k-1} \\ p_{i,j,k} \\ p_{i,j,k+1} \\ \vdots \\ p_{i,j+1,k} \\ \vdots \\ p_{i+1,j,k} \end{vmatrix} = \begin{bmatrix} \vdots \\ \nabla \cdot \vec{u} \\ \vdots \\ p_{i+1,j,k} \\ \vdots \\ p_{i+1,j,k} \\ \vdots \\ p_{i+1,j,k} \end{vmatrix}$$

$$i,j,k$$

$$Ap = d$$

- A is a huge matrix. In a NxNxN grid, A is a N³xN³ matrix.
 - I00xI00xI00 grid results in a matrix with I0¹² elements.

• A is sparse.

Boundary Conditions

 $\frac{\Delta t}{\rho} \begin{pmatrix} \frac{6p_{i,j,k} - p_{i+1,j,k} - p_{i,j+1,k} - p_{i,j,k+1}}{-p_{i-1,j,k} - p_{i,j-1,k} - p_{i,j,k-1}} \\ \frac{-p_{i-1,j,k} - p_{i,j-1,k} - p_{i,j,k-1}}{\Delta x^2} \end{pmatrix} = -\left(\frac{\frac{u_{i+1/2,j,k} - u_{i-1/2,j,k}}{\Delta x} + \frac{v_{i,j+1/2,k} - v_{i,j-1/2,k}}{\Delta x}}{+\frac{w_{i,j,k+1/2} - w_{i,j,k-1/2}}{\Delta x}} \right)$

- At cell(i,j,k), the pressures from its 6 neighboring cells are needed. What if the its neighboring cell is not Fluid?
- > At boundary cells, some modifications on the linear equation are needed.
- Each cell is either Fluid, Solid, or Empty. Since water is moving, thus the property of a cell may change (From a Empty to Fluid, or opposite) during simulating.



Empty Cell

- Assume the pressure of the Empty Cell is zero
- For example, if cell(i+1, j, k) is empty, then the linear equation at cell(i, j, k) should be:





Solid Cell

The assumption

The water do not penetrate the solid, thus

$$\vec{u} \cdot \hat{n} = \vec{u}_{\text{solid}} \cdot \hat{n}$$

If right neighboring cell (i+1,j,k) of cell(i, j,k) is Solid,



Modification on Laplace Equation

$$p_{i+1,j} = p_{i,j} + \frac{\rho \Delta x}{\Delta t} \left(u_{i+1/2,j} - u_{\text{solid}} \right)$$





Summary

$$\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} + \frac{1}{\rho} \nabla p = \vec{g}$$
$$\prod_{f_1(\vec{u})}^{"} f_3(\vec{u}) = f_2(\vec{u})$$

- $\frac{D\vec{u}}{Dt} = 0$ Advection/Transportation
- $\blacktriangleright \qquad \frac{\partial \vec{u}}{\partial t} = \vec{g}$
- **Body Forces**

Pressure/Incompressiblity

s.t. $\nabla \cdot \vec{u} = 0$



 $\vec{u}^* = \vec{u}^{advected} + \Delta t \vec{g}$

 $\mathbf{A}p = d$ $\vec{u}^{n+1} = \vec{u}^* - \frac{\Delta t}{\rho} \nabla p$

Where is the Water?

Marker particles

- initially each Fluid cell in the MAC grid will be allocated N particles (e.g. N = 4).
- Move particles in the incompressible velocity field
- update cell properties (the cell contains any particles is marked as Fluid).

From particles to Water surface

- Implicit function: f(x) = distance to the nearest particle -r
- Sample f(x) with a high resolution grid.
- Marching Cube to find Iso-surface

Water and Level Sets

Represent the surface using an implicit function

 $\{\vec{x} \mid \phi(\vec{x}) = 0\}$

- One popular choice: Signed Distance Function
 - Distance to the nearest point on the surface
 - Positive outside, negative inside.
 - Some nice properties:

$$\nabla \phi \cdot \hat{n} = 1$$
$$\hat{n} = \nabla \phi$$

Evolution of this function: advection

$$\frac{D\phi}{Dt} = 0$$

Above properties may not be preserved. Periodically recalculate the distance function.

Reference

- More details on Level Sets:
 - Book: "Level Set Methods and Dynamic Implicit Surface" by Stanley Osher and Ronald Fedkiw.
- A nice overview on Fluid Simulation
 - SIGGRAPH 2007 Course Notes, "Fluid Simulation" by Robert Bridson and Matthias Muller-Fischer.
- Libraries to Solve Sparse Linear System
 - SparseLib: <u>http://math.nist.gov/sparselib++/</u>
 - > PARDISO or IML (Intel Mathematic Library).
- Surface Reconstruction
 - Marching Cube,

http://local.wasp.uwa.edu.au/~pbourke/geometry/polygonise/

Poisson Surface Reconstruction, <u>http://www.cs.jhu.edu/~misha/</u>

Thanks!