Check the latex file and find out how macros are used to display basic quantum objects!

Quantum States. The state space \mathcal{A} of *m*-qubit is the complex Euclidean space \mathbb{C}^{2^m} . An *m*-qubit quantum state is represented by a density operator ρ , i.e., a positive semidefinite operator over \mathcal{A} with trace 1. The set of all quantum states in \mathcal{A} is denoted by Dens (\mathcal{A}).

Let $L(\mathcal{A})$ denote the set of all linear operators on space \mathcal{A} . The Hilbert-Schmidt inner product on $L(\mathcal{A})$ is defined by $\langle X, Y \rangle = \operatorname{tr}(X^*Y)$, for all $X, Y \in L(\mathcal{A})$, where X^* is the adjoint conjugate of X. Let $\operatorname{id}_{\mathcal{X}}$ denote the identity operator over \mathcal{X} , which might be omitted from the subscript if it is clear in the context. An operator $U \in L(\mathcal{X})$ is a unitary if $UU^* = U^*U = id_{\mathcal{X}}$. The set unitary operations over \mathcal{X} is denoted by $U(\mathcal{X}).$

For a multi-partite state, e.g. $\rho_{ABC} \in \text{Dens}(\mathcal{A} \otimes \mathcal{B} \otimes \mathcal{C})$, its reduced state on some subsystem(s) is represented by the same state with the corresponding subscript(s). For example, the reduced state on \mathcal{A} system of ρ_{ABC} is $\rho_A = \operatorname{tr}_{\mathcal{BC}}(\rho_{ABC})$, and $\rho_{AB} = \operatorname{tr}_{\mathcal{C}}(\rho_{ABC})$. When all subscript letters are omitted, the notation represents the original state (e.g., $\rho = \rho_{ABE}$).

A classical-quantum-, or cq-state $\rho \in \text{Dens}(\mathcal{A} \otimes \mathcal{B})$ indicates that the \mathcal{A} subsystem is classical and \mathcal{B} is quantum. Likewise for ccq-, etc., states. We use lower case letters to denote specific values assignment to the classical part of a state. For example, any cq-state $\rho_{AB} = \sum_{a} p_a |a\rangle \langle a| \otimes \rho_B^a$ in which $p_a = \mathbf{Pr}[A = a]$ and ρ_B^a is a normalized state.

Distance Measures. For any $X \in L(\mathcal{A})$ with singular values $\sigma_1, \dots, \sigma_d$, where $d = \dim(\mathcal{A})$, the trace norm of \mathcal{A} is $||X||_{\text{tr}} = \sum_{i=1}^{d} \sigma_i$. The *trace distance* between two quantum states ρ_0 and ρ_1 is defined to be

$$|\rho_0 - \rho_1|_{\rm tr} \stackrel{\rm def}{=} \frac{1}{2} \|\rho_0 - \rho_1\|_{\rm tr}.$$

, which admits the following operational meaning. The following well-known fact relates the trace distance with the optimal probability of distinguishing quantum states. Their *fidelity*, denoted by $F(\rho_0, \rho_1)$), is

$$F(\rho_0, \rho_1) = \|\sqrt{\rho_0}\sqrt{\rho_1}\|_{tr}.$$
 (1)

When ρ_0 and ρ_1 are *classical* states, the trace distance $|\rho_0 - \rho_1|_{tr}$ is equivalent to the *statistical* distance between ρ_0 and ρ_1 . It is also a well known fact that for two distributions X_1, X_2 over \mathcal{X} , let $p_x = \Pr[X_1 = x]$ and $q_x = \mathbf{Pr}[X_2 = x]$ and their statistical distance satisfies

$$|X_1 - X_2|_{\rm tr} = \frac{1}{2} \sum_x |p_x - q_x| = \sum_{x: p_x > q_x} (p_x - q_x).$$
(2)

For simplicity, when both states are classical, we use $(X_1) \approx_{\epsilon} (X_2)$ to denote $|X_1 - X_2|_{tr} \leq \epsilon$.

Moreover, the trace distance admits the following two simple facts.

Fact 1. For any state $\rho_1, \rho_2 \in \text{Dens}(\mathcal{A})$ and $\sigma \in \text{Dens}(\mathcal{B})$, we have

$$\left|\rho_{1}-\rho_{2}\right|_{\mathrm{tr}}=\left|\rho_{1}\otimes\sigma-\rho_{2}\otimes\sigma\right|_{\mathrm{tr}}$$

Fact 2. Let $\rho, \sigma \in \text{Dens}(\mathcal{A} \otimes \mathcal{B})$ be any two cq-states where \mathcal{A} is the classical part. Moreover, $\rho =$ $\sum_{a} p_a |a\rangle \langle a| \otimes \rho_B^a$ and $\sigma = \sum_{a} q_a |a\rangle \langle a| \otimes \sigma_B^a$. Then we have

$$\left|\rho - \sigma\right|_{\mathrm{tr}} = \sum_{a} \left|p_{a}\rho_{B}^{a} - q_{a}\sigma_{B}^{a}\right|_{\mathrm{tr}}.$$

References