# Quantum walks on graphs

#### Andrew Childs California Institute of Technology



# Information is physical.

## Outline

- I. Introduction: Quantum systems as information processing machines
- 2. Exponential speedup by quantum walk
- 3. Spatial search by quantum walk
- 4. Evaluating Boolean formulas

## I. Introduction

## The universal quantum computer

(The ultimate quantum physics lab!)

- ullet Prepare n qubits in the state  $|00\dots0
  angle$
- Apply a sequence of unitary operations acting on one or two qubits at a time
- Perform a measurement to get the result



Note: Many equivalent models exist (Hamiltonian dynamics of coupled spins, braiding of nonabelian anyons, quantum cellular automata, ...).

## Implementations of quantum computers

#### Trapped ions



(Monroe & Wineland)

#### Quantum dots



Nuclear spins



(Chuang et al.)

#### Superconducting circuits



(Nakamura et al.)

... and many others!

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In the rest of this talk, we assume a perfectly functioning quantum computer.

Toy problem [Deutsch 1985]

Given a black box for  $f \colon \{0,1\} \to \{0,1\}$ , determine  $f(0) \oplus f(1)$ .

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Classically: Two queries required.

Quantumly: One query sufficient!



But don't classical systems exhibit interference too?

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How can we exploit the efficient representation of interference phenomena to perform fast computations?





$$A = \begin{pmatrix} 0 & 1 & 1 & 0 & 0 \\ 1 & 0 & 0 & 1 & 1 \\ 1 & 0 & 0 & 1 & 0 \\ 0 & 1 & 1 & 0 & 1 \\ 0 & 1 & 0 & 1 & 0 \end{pmatrix}$$

adjacency matrix



		/0	1	1	0	0			(-2)
		1	0	0	1	1			1
A	=	1	0	0	1	0		L =	1
		0	1	1	0	1			0
		$\left( 0 \right)$	1	0	1	0/			$\int 0$
adjacency matrix							ix		Lapl

$$\begin{pmatrix} -2 & 1 & 1 & 0 & 0 \\ 1 & -3 & 0 & 1 & 1 \\ 1 & 0 & -2 & 1 & 0 \\ 0 & 1 & 1 & -3 & 1 \\ 0 & 1 & 0 & 1 & -2 \end{pmatrix}$$

1

Laplacian



Random walk on G

State: Probability  $p_j(t)$  of being at vertex j at time tDynamics:  $\frac{d}{dt}\vec{p} = -\gamma L\vec{p}$ 



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State: Probability  $p_j(t)$  of being at vertex j at time tDynamics:  $\frac{d}{dt}\vec{p} = -\gamma L\vec{p}$ 

Quantum walk on G

State: Amplitude  $q_j(t)$  to be at vertex j at time tDynamics:  $i \frac{d}{dt} \vec{q} = -\gamma L \vec{q}$ 

# 2. Exponential speedup by quantum walk

• Childs, Farhi, and Gutmann, Quantum Inf. Proc. 1, 35-43 (2002).

• Childs, Cleve, Deotto, Farhi, Gutmann, and Spielman, in Proc. STOC (2003), pp. 59-68.



## Black box graph traversal problem

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Claim: There is a family of graphs  $G_n$  (with designated in and out vertices) for which a quantum walk starting at the in vertex finds the out vertex in time poly(n), but any classical algorithm using poly(n) queries finds the out vertex with exponentially small probability.

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## Reduction of the quantum walk



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#### Column subspace

$$|\text{col } j\rangle = \frac{1}{\sqrt{N_j}} \sum_{a \in \text{column } j} |a\rangle$$

where

$$N_{j} = \begin{cases} 2^{j} & 0 \le j \le n \\ 2^{2n+1-j} & n+1 \le j \le 2n+1 \end{cases}$$
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Reduced adjacency matrix

$$\operatorname{col} j|A|\operatorname{col} j+1\rangle = \begin{cases} \sqrt{2} & 0 \leq j \leq n-1, \\ & n+1 \leq j \leq 2n \\ 2 & j=n \end{cases}$$

Problem: Given a black box for G, implement the quantum walk on G, i.e., simulate the unitary time evolution  $e^{-iHt}$  where H=L (or A). (Cf. implementing a random walk, which is easy.)

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Main idea: Color the graph. Then the simulation breaks into small pieces that are easy to handle.



Any sufficiently sparse graph can be efficiently colored using only local information (Linial 1987).

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This idea is useful for simulating Hamiltonians whenever the graph of nonzero matrix elements is sufficiently sparse.

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• Childs and Goldstone, Phys. Rev. A 70, 022314 (2004).

• Childs and Goldstone, Phys. Rev. A 70, 042312 (2004).

Search space: N items,  $x \in \{1, 2, ..., N\}$ Goal: Find one marked item, w Query: "is w = x?" i.e., black box function  $f(x) = \begin{cases} 0 & x \neq w \\ 1 & x = w \end{cases}$ 

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Grover searching can be applied to many computational problems. But can it be used to search a physical database, in which the N items are distributed in space?

Quantum walk approach:  $H = -\gamma L - |w\rangle \langle w|$ 

$$L_{ab} = \begin{cases} 1 & ab \in G \\ -\deg(a) & a = b \\ 0 & \text{otherwise} \end{cases}$$
$$L \approx \nabla^2$$

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$$\gamma o 0$$
  
 $H pprox - |w\rangle \langle w|$   
ground state  $pprox |w\rangle$   
first excited state  $pprox |s\rangle$ 

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## Results

Graph	Success amplitude	Run time
Complete	1 - o(1)	$O(N^{1/2})$
Hypercube	1 - o(1)	$O(N^{1/2})$
Lattice, $d > 4$	O(1)	$O(N^{1/2})$
Lattice, $d = 4$	$O(1/{\log^{1/2}N})$	$O((N \log N)^{1/2})$
Lattice, $d = 3$	$O(N^{\text{-}1/6})$	$O(N^{2/3})$
Lattice, $d = 2$	$O((\log N/N)^{1/2})$	$O(N/{\log N})$

### Behavior for $d{>}4$

$$\begin{array}{ll} \text{Critical } \gamma: & \gamma_* = I_{1,d} \\ \text{Run time:} & T = \frac{\pi \sqrt{I_{2,d}N}}{2I_{1,d}} \\ \text{Success probability:} & |\langle w|e^{-iHT}|s\rangle|^2 = \frac{I_{1,d}^2}{I_{2,d}} \end{array}$$

#### where

$$I_{j,d} = \frac{1}{(2\pi)^d} \int_{-\pi}^{\pi} \frac{\mathrm{d}^d \vec{k}}{[\mathcal{E}(\vec{k})]^j}$$
$$\mathcal{E}(\vec{k}) = 2\gamma \left(d - \sum_{j=1}^d \cos k_j\right)$$

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Note: 
$$\int_{0} \frac{\mathrm{d}^{d}\vec{k}}{|\vec{k}|^{p}} \sim \int_{0} \frac{k^{d-1}\mathrm{d}k}{|\vec{k}|^{p}}$$
converges for  $d > p$ .

#### dispersion relation



# The Dirac equation: Faster search for d = 2, 3, 4Dirac Hamiltonian: $H_{\text{Dirac}} = \sum_{j=1}^{d} \alpha_j p_j + \beta m$ , $p_j = i \frac{d}{dx_j}$ with $\{\alpha_j, \alpha_k\} = 2\delta_{jk}$ , $\{\alpha_j, \beta\} = 0$ , $\beta^2 = 1$

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Lattice version: 
$$H_0 = \omega \sum_{j=1}^d \alpha_j P_j$$
,  $P_j |\vec{x}\rangle = \frac{i}{2} (|\vec{x} + \hat{e}_j\rangle - |\vec{x} - \hat{e}_j\rangle)$   
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Improved:  $H_0 = \omega \sum_{j=1}^d \alpha_j P_j + \gamma \beta L$  $\mathcal{E}(\vec{k}) = \pm \sqrt{\omega^2 \sum_{j=1}^d \sin^2 k_j} + \gamma^2 [2 \sum_{j=1}^d (1 - \cos k_j)]^2$ 

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#### Algorithm:

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#### Run time:

 $\begin{array}{l} O(\sqrt{N}) \text{ in } d \! > \! 2\text{,} \\ O(\sqrt{N}\log N) \text{ in } d \! = \! 2\text{.} \end{array}$ 



• Childs, Cleve, Jordan, and Yeung, quant-ph/0702160.

• Childs, Reichardt, Špalek, and Zhang, quant-ph/0703015.

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Given a black box for determining how the variables are assigned, how many variables must we query to determine the value of the formula?

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Example: Unstructured search (OR)





Classical complexity:  $\Theta(N)$ Quantum algorithm [Grover 1996]:  $O(\sqrt{N})$ Quantum lower bound [BBBV 1996]:  $\Omega(\sqrt{N})$ 















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Discretized version [Childs-Cleve-Jordan-Yeung 07]:  $O(\sqrt{N^{1+\epsilon}})$ 










### Evaluating AND-OR trees by scattering



Claim: For small k, the transmission coefficient is large if the formula (translated into NAND gates) evaluates to 0, and small if it evaluates to 1.





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Main idea: The quantum walk on the (expanded, appropriately weighted) tree has a zero eigenvalue if the formula evaluates to 1, and a smallest eigenvalue  $O(\frac{1}{\sqrt{N}})$  if it evaluates to 0. Use phase estimation.

# Summary and outlook

- Physics ↔ Information
- Quantum systems can encode and process information in a fundamentally non-classical way.
- While we have many examples of this phenomenon, we are far from having a general understanding of the information processing power of quantum mechanics.
- In particular, we would like to develop new ways of exploiting quantum interference to perform information processing tasks.