A Verified Optimizer for Quantum Circuits

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We present voqc, the first fully verified compiler for quantum circuits, written using the Coq proof assistant. Quantum circuits are expressed as programs in a simple, low-level language called sqir, which is deeply embedded in Coq. Optimizations and other transformations are expressed as Coq functions, which are proved correct with respect to a semantics of sqir programs. sqir uses a semantics of matrices of complex numbers, which is the standard for quantum computation, but treats matrices symbolically in order to reason about programs that use an arbitrary number of quantum bits. sqir’s careful design and our provided automation make it easy to write and verify a broad range of optimizations in voqc, and even allow us to verify correctness properties of interesting source sqir programs.

CCS Concepts: •Hardware → Quantum computation; Circuit optimization; •Software and its engineering → Formal software verification;

Additional Key Words and Phrases: Formal Verification, Quantum Computing, Circuit Optimization, Certified Compilation, Programming Languages

1 INTRODUCTION

Programming quantum computers will be challenging, at least in the near term. Qubits will be scarce, and gate pipelines will need to be short to prevent decoherence. Fortunately, optimizing compilers can transform a source algorithm to work with fewer resources. Where compilers fall short, programmers can optimize their algorithms by hand.

Of course, both compiler and by-hand optimizations will inevitably have bugs. As evidence of the former, Kissinger and van de Wetering [2019] discovered mistakes in the optimized outputs produced by the circuit compiler by Nam et al. [2018], and Nam et al. themselves found that the optimization library they compared against (Amy et al. [2013]) sometimes produced incorrect results. Likewise, Amy [2018] discovered an optimizer they had recently developed produced buggy results [Amy et al. 2018]. Making mistakes when optimizing by hand is also to be expected: as put well by Zamdzhiev [2016], quantum computing can be frustratingly unintuitive.

Unfortunately, the very factors that motivate optimizing quantum compilers make it difficult to test their correctness. Comparing runs of a source program to those of its optimized version is often impractical due to the indeterminacy of typical quantum algorithms and the substantial expense involved in executing or simulating them. Indeed, resources may be too scarce, or the qubit connectivity too constrained, to run the program without optimization!

An appealing solution to this problem is to apply rigorous formal methods to prove that an optimization or algorithm always does what it is intended to do. For example, CompCert [Leroy 2009] is a compiler for C programs that is written and proved correct using the Coq proof assistant [Coq Development Team 2019]. CompCert includes sophisticated optimizations whose proofs of correctness are verified to be valid by Coq’s type checker.

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In this paper, we apply CompCert’s approach to the quantum setting. We present \texttt{voqc} (pronounced “vox”), a verified optimizer for quantum circuits. \texttt{voqc} takes as input a quantum program written in a language we call \texttt{sqir} (“squire”). \texttt{sqir} is designed to be a simple quantum intermediate representation, but it is suitable for source-level programming too: it is not very different from languages such as Quil [Smith et al. 2016] or OpenQASM [Cross et al. 2017], which describe quantum programs as circuits. \texttt{sqir} is deeply embedded in Coq, similar to how Quil is embedded in Python via PyQuil [Rigetti Computing 2019a], allowing us to write sophisticated quantum programs. \texttt{voqc} applies a series of optimizations to \texttt{sqir} programs, ultimately producing a result that is compatible with the specified quantum architecture. For added convenience, \texttt{voqc} provides translators between \texttt{sqir} and OpenQASM. (Section 2.)

We designed \texttt{sqir} to make it as easy as possible to reason about the semantics of quantum programs, which are significantly different from the semantics of classical programs. For example, while classical computing allows us to reason about different variables in a program independently, the phenomenon of quantum entanglement requires us to think about a global quantum state, typically represented as a large vector or matrix of complex numbers. \texttt{sqir} uses natural numbers in place of variables so that we can naturally index into the vector or matrix state. \texttt{sqir} moreover provides two semantics for quantum programs: For circuits without measurement, we use a unitary semantics, where every program corresponds to a restricted class of square matrices. For those programs that do employ measurement, we express their semantics as a function between density matrices. These are harder to work with but can represent the broader class of mixed states. Both semantics have proved widely useful: In addition to verifying program transformations, we use \texttt{sqir} to directly prove the correctness of several quantum programs, including GHZ state preparation, quantum teleportation, superdense coding, the Deutsch-Jozsa algorithm, and quantum phase estimation (Section 3).

At the core of \texttt{voqc} is a framework for writing transformations of \texttt{sqir} programs and verifying their correctness. To ensure that the framework is suitably expressive, we have used it to develop verified versions of a variety of optimizations. Many are based on those used in an optimizer developed by Nam et al. [2018], which is the best performing optimizer we know of (per experiments described below). We abstract these optimizations into a couple of different classes, and provide library functions, lemmas, and automation to simplify their construction and proof. We have also verified a circuit mapping routine that transforms \texttt{sqir} programs to satisfy constraints on how qubits may interact on a particular target architecture. All of these transformations were reasonably straightforward to prove correct thanks to \texttt{sqir}’s design. (Sections 4 and 5.)

We evaluated the quality of the optimizations we verified in \texttt{voqc}, and by extension the quality of its framework, by measuring how well it optimizes a set of benchmark programs, compared to Nam et al. and several other compilers. The results are encouraging. On a benchmark of 28 circuit programs developed by Amy et al. [2013] we find that \texttt{voqc} reduces total gate count on average by 18.7% compared to 10.4% by IBM’s Qiskit compiler [Aleksandrowicz et al. 2019], 10.9% by CQC’s tket compiler [Cambridge Quantum Computing Ltd 2019], and 26.4% by Nam et al. On the same benchmarks, \texttt{voqc} reduces $T$-gate count (an important measure when considering fault tolerance) on average by 42.3% compared to 40.9% by Amy et al., 42.3% by Nam et al., and 43.8% by the PyZX compiler [Kissinger and van de Wetering 2019]. Results on an even larger benchmark suite (detailed in Appendix C) tell the same story. In sum, \texttt{voqc} and \texttt{sqir} are expressive enough to verify a range of useful optimizations, yielding performance competitive with standard compilers. (Section 6.)

\texttt{voqc} is the first fully verified optimizer for general quantum programs. Amy et al. [2017] developed a verified optimizing compiler from source Boolean expressions to reversible circuits and Fagan and Duncan [2018] verified an optimizer for ZX-diagrams representing Clifford circuits; however, neither of these tools handle general quantum programs. In concurrent work, Shi et al.
[2019] developed CertiQ, which uses symbolic execution and SMT solving to verify some circuit transformations in the Qiskit compiler. CertiQ is limited to verifying correct application of local equivalences and does not provide a way to describe general quantum state (a key feature of sQIR), which limits the types of optimizations that it can reason about. This also means that it cannot be used as a tool for verifying general quantum programs. Smith and Thornton [2019] presented a compiler with built-in translation validation via QMDD equivalence checking [Miller and Thornton 2006]. However, QMDDs represent quantum state concretely, which means that the validation time will increase exponentially with the number of qubits in the compiled program. In contrast to these, sQIR represents matrices symbolically, which allows us to reason about arbitrary quantum computation and verify interesting, non-local optimizations, independently of the number of qubits in the optimized program. (Section 7.)

Our work on voqc and sQIR are steps toward a broader goal of developing a full-scale verified compiler toolchain. Next steps include developing certified transformations from higher-level quantum languages to sQIR and implementing optimizations with different objectives, e.g., that aim to reduce the probability that a result is corrupted by quantum noise. All code we reference in this paper can be found online at https://github.com/inQWIRE/SQIR.

2 OVERVIEW

We begin with a brief background on quantum programs, focusing on the challenges related to formal verification. We then provide an overview of voqc and sQIR, summarizing how they address these challenges.

2.1 Preliminaries

Quantum programs operate over quantum states, which consist of one or more quantum bits (aka, qubits). A single qubit is represented as a vector of complex numbers $\langle \alpha, \beta \rangle$ such that $|\alpha|^2 + |\beta|^2 = 1$. The vector $\langle 1, 0 \rangle$ represents the state $|0\rangle$ while vector $\langle 0, 1 \rangle$ represents the state $|1\rangle$. A state written $|\psi\rangle$ is called a ket, following Dirac’s notation. We say a qubit is in a superposition of $|0\rangle$ and $|1\rangle$ when both $\alpha$ and $\beta$ are non-zero. Just as Schrodinger’s cat is both dead and alive until the box is opened, a qubit is only in superposition until it is measured, at which point the outcome will be 0 with probability $|\alpha|^2$ and 1 with probability $|\beta|^2$. Measurement is not passive: it has the effect of collapsing the state to match the measured outcome, i.e., either $|0\rangle$ or $|1\rangle$. As a result, all subsequent measurements return the same answer.

Operators on quantum states are linear mappings. These mappings can be expressed as matrices, and their application to a state expressed as matrix multiplication. For example, the Hadamard operator $H$ is expressed as a matrix $\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$. Applying $H$ to state $|0\rangle$ yields state $\left( \frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}} \right)$, also written as $|+\rangle$. Many quantum operators are not only linear, they are also unitary—the conjugate transpose (or adjoint) of their matrix is its own inverse. This ensures that multiplying a qubit by the operator preserves the qubit’s sum of norms squared. Since a Hadamard is its own adjoint, it is also its own inverse: hence $H |+\rangle = |0\rangle$.

A quantum state with $n$ qubits is represented as vector of length $2^n$. For example, a 2-qubit state is represented as a vector $\langle \alpha, \beta, \gamma, \delta \rangle$ where each component corresponds to (the square root of) the probability of measuring $|00\rangle$, $|01\rangle$, $|10\rangle$, and $|11\rangle$, respectively. Because of the exponential size of the complex quantum state space, it is not possible to simulate a 100-qubit quantum computer using even the most powerful classical computer!

$n$-qubit operators are represented as $2^n \times 2^n$ matrices. For example, the CNOT operator over two qubits is expressed as the matrix shown at the right. It expresses a controlled not
operation—if the first qubit (called the *control*) is $|0\rangle$ then both qubits are mapped to themselves, but if the first qubit is $|1\rangle$ then the second qubit (called the *target*) is negated, e.g., $\text{CNOT}|00\rangle = |00\rangle$ while $\text{CNOT}|10\rangle = |11\rangle$.

$n$-qubit operators can be used to create *entanglement*, which is a situation where two qubits cannot be described independently. For example, while the vector $\langle 1, 0, 0, 0 \rangle$ can be written as $\langle 1, 0 \rangle \otimes \langle 1, 0 \rangle$ where $\otimes$ is the tensor product, the state $\langle \frac{1}{\sqrt{2}}, 0, 0, \frac{1}{\sqrt{2}} \rangle$ cannot be similarly decomposed. We say that $\langle \frac{1}{\sqrt{2}}, 0, 0, \frac{1}{\sqrt{2}} \rangle$ is an entangled state.

An important non-unitary quantum operator is *projection* onto a subspace. For example, $|0\rangle\langle 0|$ (in matrix notation $\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$) projects a qubit onto the subspace where that qubit is in the $|0\rangle$ state. Projections are useful for describing quantum states after measurement has been performed. We sometimes use $\langle i \rangle_q \langle i |$ as shorthand for applying the projection $|i\rangle \langle i|$ to qubit $q$ and an identity operation to every other qubit in the state.

### 2.2 Quantum Circuits

Quantum programs are typically expressed as circuits, as shown in Figure 1(a). In these circuits, each horizontal wire represents a *qubit* and boxes on these wires indicate quantum operators, or *gates*. Gates can either be unitary operators (e.g., Hadamard, CNOT) or non-unitary ones (e.g., measurement). In software, quantum circuit programs are often represented using lists of instructions that describe the different gate applications. For example, Figure 1(b) is the Quil [Smith et al. 2016] representation of the circuit in Figure 1(a).

In the *QRAM model* [Knill 1996] quantum computers are used as co-processors to classical computers. The classical computer generates descriptions of circuits to send to the quantum computer and then processes the measurement results. High-level quantum programming languages are designed to follow this model. For example, Figure 1(c) shows a program in PyQuil [Rigetti Computing 2019a], a quantum programming framework embedded in Python. The *ghz_state* function takes an array *qubits* and constructs a circuit that prepares the Greenberger-Horne-Zeilinger (GHZ) state [Greenberger et al. 1989], which is an $n$-qubit entangled quantum state of the form

$$|\text{GHZ}^n\rangle = \frac{1}{\sqrt{2}}(|0\rangle^\otimes n + |1\rangle^\otimes n).$$

Calling *ghz_state([0, 1, 2])* would return the Quil program in Figure 1(b), which the quantum computer could subsequently execute. The high-level language may provide facilities to optimize constructed circuits, e.g., to reduce gate count, circuit depth, and qubit usage. It may also perform transformations to account for hardware-specific details like the number of qubits, available set of gates, or connectivity between physical qubits.
What if we want to formally verify that ghz_state, when passed an array of indices \([0, \ldots, n-1]\), returns a circuit that produces the quantum state \(|\text{GHZ}^n\rangle\)? What steps are necessary?

First, we need a way to formally define quantum states as matrices of complex numbers. Indeed, we need a way to define indexed families of states—\(|\text{GHZ}^n\rangle\) is a function from an index \(n\) to a quantum state. Second, we need a formal language in which to express quantum programs; to this language we must ascribe a mathematical semantics in terms of quantum states. A program like ghz_state is a function from an index (a list of length \(n\)) to a circuit (of size \(n\)), and this circuit’s denotation is its equivalent (unitary) matrix (of size \(2^n\)). Finally, we need a way to mechanically reason that, for arbitrary \(n\), the semantics of ghz_state([0, 1, \ldots, n−1]) applied to the zero state \((|0\rangle^\otimes n)\) is equal to the state \(|\text{GHZ}^n\rangle\).

We designed sqir, a simple quantum intermediate representation, to do all of these things. sqir is a simple circuit-oriented language deeply embedded in the Coq proof assistant in a manner similar to how Quil is embedded in Python via PyQuil. We use sqir’s host language, Coq, to define the syntax and semantics of sqir programs, and to express properties involving quantum states. We developed a library of lemmas and tactic-based automation to assist in writing proofs about quantum programs; such proofs make heavy use of complex numbers and linear algebra. These proofs are aided by isolating sqir’s unitary core from primitives for measurement, which require consideration of probability distributions of outcomes (represented as density matrices); this means that (sub-)programs that lack measurement can have simpler proofs. Either way, in sqir we perform reasoning symbolically. For example, we can prove that every circuit generated by the sqir-equivalent of ghz_state produces the corresponding state \(|\text{GHZ}^n\rangle\) when applied to input lists of length \(n\), for any \(n\).

sqir is implemented in just over 3300 lines of Coq, with an additional 2000 lines of example sqir programs and proofs. We started with Coq libraries for complex numbers and matrices developed for the Qwire language [Paykin et al. 2017]; over the course of our work we have extended these libraries with around 3000 lines of code providing more automation for linear algebra and better support for complex phases. We present sqir’s syntax and semantics along with example programs and verified properties of correctness in Section 3.

2.4 voqc: A Verified Optimizer for Quantum Circuits

While sqir is suitable for proving correctness properties about source programs (like ghz_state), its primary use has been as the intermediate representation of voqc, our verified optimizer for quantum circuits, and the signature achievement of this paper. An optimizer is a function from programs to programs, with the intention that the output program has the same semantics as the input. In voqc, we prove this is always the case: a voqc optimization \(f\) is a Coq function over sqir circuit \(C\), and we prove that the semantics of input circuit \(C\) is always equivalent to the semantics of the output \(f(C)\).

The voqc approach stands in contrast to prior work that relies on translation validation [Amy 2018; Kissinger and van de Wetering 2019; Smith and Thornton 2019], which may fail to identify latent bugs in the optimizer, while adding run-time overhead. By proving correctness with respect to an explicit semantics for input/output programs (i.e., that of sqir), voqc optimizations are flexible in their expression. Prior work has been limited to the application of local equivalences [Shi et al. 2019], leaving highly effective, global transformations we have proved correct in voqc out of reach. Such global proofs are aided by the design of sqir (notably, the isolation of a unitary core) and accompanying proof automation.
The structure of voqc is summarized in Figure 2. The voqc transformations themselves are shown at the middle right, and are described in Sections 4 and 5. In addition to performing circuit optimizations, voqc also performs circuit mapping, transforming a sqir program to an equivalent one that respects constraints imposed by the target architecture. Once again, we prove that it does so correctly.

Using Coq’s standard code extraction mechanism, we can extract voqc into a standalone OCaml program. This program takes as input a sqir program in an OCaml representation. This input can be extracted from a Coq-hosted (and proved correct) sqir program (upper right), or from a program expressed in OpenQASM [Cross et al. 2017], a standard representation for quantum circuits (upper left). Since a number of quantum programming frameworks, including Qiskit [Aleksandrowicz et al. 2019], t\ket [Cambridge Quantum Computing Ltd 2019], Project Q [Steiger et al. 2018] and Cirq [The Cirq Developers 2019], can output OpenQASM, this allows us to run voqc on a variety of generated circuits, without requiring the user to program in OCaml or Coq.

voqc is implemented in about 7200 lines of Coq, with roughly 2100 lines for circuit mapping, 2100 lines for general-purpose sqir program manipulation, 2200 lines for unitary program optimizations, and 800 lines for non-unitary program optimizations. We use about 300 lines of standalone OCaml code for running voqc on our benchmarks in Section 6.

3 \textbf{SQIR: A SMALL QUANTUM INTERMEDIATE REPRESENTATION}

This section presents the syntax and semantics of sqir programs. We begin with the core of sqir, which describes unitary circuits. We then describe the expanded language, which allows measurement and initialization. As illustrations of sqir’s suitability for source-program verification, we describe sqir proofs of correctness for GHZ state preparation (a unitary program) and quantum teleportation (a non-unitary program); Appendix B presents additional examples.
\[ \begin{aligned}
\begin{bmatrix} U_1; U_2 \end{bmatrix}_d &= \begin{bmatrix} U_2 \end{bmatrix}_d \times \begin{bmatrix} U_1 \end{bmatrix}_d \\
\begin{bmatrix} G_1 \ q \end{bmatrix}_d &= \begin{cases} \text{apply}_1(G_1, \ q, \ d) & \text{well-typed} \\
0_{2d} & \text{otherwise} \end{cases} \\
\begin{bmatrix} G_2 \ q_1 \ q_2 \end{bmatrix}_d &= \begin{cases} \text{apply}_2(G_2, \ q_1, \ q_2, \ d) & \text{well-typed} \\
0_{2d} & \text{otherwise} \end{cases}
\end{aligned} \]

Fig. 3. Semantics of unitary sqir programs, assuming a global register of dimension \( d \). The \text{apply}_k function maps a gate name to its corresponding unitary matrix and extends the intended operation to the given dimension by applying an identity operation on every other qubit in the system.

3.1 Unitary Core

The sqir language is composed of two parts: a core language of unitary operators, and a full language that adds measurement to this core. This design eases formal proof. As we show here, the semantics of a unitary sqir program is expressed directly as a matrix, which means that proofs of correctness of unitary optimizations (the bulk of voqc) involve reasoning directly about matrices. Doing so is far simpler than reasoning about functions over density matrices, as is required for a language (like full sqir) involving measurement.

3.1.1 Syntax. A unitary sqir program \( U \) is a sequence of applications of gates \( G \) to qubits \( q \).

\[ U := U_1; U_2 | G \ q | G \ q_1 \ q_2 \]

Qubits are referred to by natural numbers that index into a global register of quantum bits. Each sqir program is parameterized by a set of unitary one- and two-qubit gates (from which \( G \) is drawn) and the dimension of the global register (i.e., the number of available qubits). In Coq, a unitary sqir program \( U \) has type \( \text{ucom} \ n \), where \( n \) identifies the gate set and \( n \) is the size of the global register.

As an example, consider the program to the right, which is equivalent to PyQuil’s \text{ghz}_\text{state} from Figure 1(c). The Coq function \text{ghz} recursively constructs a sqir program, i.e., a Coq value of type \( \text{ucom} \ n \). This program, when run, prepares the GHZ state. When \( n = 0 \), \text{ghz} produces a sqir program that is just the identity gate \( I \) applied to qubit 0. When \( n = 1 \), the result is the Hadamard gate \( H \) applied to qubit 0. When \( n > 1 \), \text{ghz} constructs the program \( U_1; U_2 \), where \( U_1 \) is the \text{ghz} circuit on \( n - 1 \) qubits, and \( U_2 \) is the appropriate CNOT gate. The result of \text{ghz} 3 is equivalent to the circuit shown in Figure 1(a).

3.1.2 Semantics. Suppose that \( u_1 \) and \( u_2 \) are the matrices corresponding to unitary programs \( U_1 \) and \( U_2 \) and \( |\psi\rangle \) is a quantum state vector. Recall that matrix multiplication is associative, so \( u_2 (u_1 |\psi\rangle) \) is equivalent to \((u_2u_1) |\psi\rangle\). Moreover, recall that multiplying two unitary matrices yields a unitary matrix. As such, the semantics of sqir program \( U_1; U_2 \) is naturally described by the unitary matrix \( u_2u_1 \).

This semantics is shown in Figure 3. There are two things to notice. First, if a program is not well-typed its denotation is the zero matrix (of size \( 2^d \times 2^d \)). A program \( U \) is well-typed if every gate application is valid, meaning that its index arguments are within the bounds of the global register, and no index is repeated. The latter requirement enforces linearity and thereby quantum
mechanics’ no-cloning theorem, which says that it is impossible to create a copy of an arbitrary quantum state.

Otherwise, the program’s denotation follows from the composition of the matrices that correspond to each of the applications of its unitary gates, \( G \). The only wrinkle is that a full program consists of many gates operating on a subset (1 or 2) of the total qubits; thus, a gate application’s matrix needs to apply the identity operation to the other qubits. This is what \( \text{apply}_1 \) and \( \text{apply}_2 \) do. For example, \( \text{apply}_1(G_u, q, d) = I_{2q} \otimes u \otimes I_{2(d-q-1)} \) where \( u \) is the matrix interpretation of the gate \( G_u \) and \( I_k \) is the \( k \times k \) identity matrix. The \( \text{apply}_2 \) function requires us to decompose the two-qubit unitary into a sum of tensor products: for instance, \( \text{CNOT} \) can be written as \( |0\rangle \langle 0| \otimes I_2 \otimes |1\rangle \otimes \sigma_x \) where \( \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \). We then have

\[
\text{apply}_2(\text{CNOT}, q_1, q_2, d) = I_{2q_1} \otimes |0\rangle \langle 0| \otimes I_{2q_2} \otimes I_2 \otimes I_{2d-2q_1-2q_2} + I_{2q_1} \otimes |1\rangle \otimes I_{2q_2} \otimes \sigma_x \otimes I_{2d-2q_1-2q_2}
\]

where \( r = q_2 - q_1 - 1 \) and \( s = d - r - 1 \).

In our development we define the semantics of \( \text{sqir} \) programs over the gate set \( G \in \{ R_{\theta, \phi, \lambda}, \text{CNOT} \} \) where \( R_{\theta, \phi, \lambda} \) is a general single-qubit rotation parameterized by three real-valued rotation angles and \( \text{CNOT} \) is the standard two-qubit controlled-not gate. This is our base set of gates. It is the same as the underlying set used by OpenQASM [Cross et al. 2017] and is universal, meaning that it can approximate any unitary operation to within arbitrary error. The matrix interpretation of the single-qubit \( R_{\theta, \phi, \lambda} \) gate is

\[
\begin{pmatrix}
\cos(\theta/2) & -e^{i\lambda} \sin(\theta/2) \\
e^{i\phi} \sin(\theta/2) & e^{i(\phi+\lambda)} \cos(\theta/2)
\end{pmatrix}
\]

and the matrix interpretation of the \( \text{CNOT} \) gate is given in Section 2.1.

Common single-qubit gates can be defined in terms of \( R_{\theta, \phi, \lambda} \). For example, the two single-qubit gates used in our GHZ example—identity \( I \) and Hadamard \( H \)—are defined as \( R_{0,0,0} \) and \( R_{\pi/2,0,\pi} \), respectively. The Pauli \( X \) gate is \( R_{\pi/2,0,\pi} \) and the Pauli \( Z \) gate is \( R_{0,0,\pi} \). We can also define more complex operations as \( \text{sqir} \) programs. For example, the \( \text{SWAP} \) operation, which swaps two qubits, can be defined as a sequence of three \( \text{CNOT} \) gates.

3.1.3 Proofs. We designed \( \text{sqir} \)’s unitary core to simplify formal proofs, in three ways.

First, by giving a denotation of the zero matrix to ill-typed gate applications (and thereby ill-typed programs), we do not need to explicitly assume or prove that a program is well-typed in order to state a property about its semantics, thereby removing clutter from theorems and proofs. For example, in our proof of the \( \text{ghz} \) program below we do not need to explicitly prove that \( \text{ghz} \) is well-typed (although this is true).

Second, we do not use dependent types to represent matrices in the semantics. Following Rand et al. [2017, 2018a], we define matrices as functions from pairs of natural numbers to complex numbers.

**Definition** Matrix \((m : \mathbb{N}) : \mathbb{N} \rightarrow \mathbb{N} \rightarrow \mathbb{C} \).

The arguments \( m \) and \( n \), which are the dimensions of the matrix, are _phantom types_—they do not appear in the definition. The phantom types are useful to define certain operations on matrices that depend on these dimensions, such as the Kronecker product and matrix multiplication. However, there is no proof burden internal to the matrices themselves. Instead, it is possible to show a matrix is well-formed within its specified bounds by means of an external predicate:

**Definition** WF_Matrix \((m : \text{Matrix} m n) : \mathbb{P} := \forall i, j, i \geq m \vee j \geq n \rightarrow M i j = 0 \).

 Phantom types occupy a convenient middle ground in allowing information to be stored in the types, while pushing the majority of the work to external predicates. For instance, we define \(|i|^{\otimes n}\) (used below) recursively as \(|i| \otimes |i| \otimes \cdots \otimes |i|\), with \( n \) repetitions. Coq has no way of inferring a type
for this, so we declare that is has type \( \text{Vector} \ 2^n \), and Coq will allow us to use it in any context where a vector is expected. However, before using it in rewrite rules like \( I_2^n \times |j\rangle^{\otimes n} = |j\rangle^{\otimes n} \) (which says that multiplication by the identity matrix is an identity operation), we will need to show that \(|j\rangle^{\otimes n}\) is a well-formed \(2^n\) length vector.

Lastly, and most importantly, the semantics of unitary sqir is straightforward thanks to its use of concrete (numeric) indices to refer to qubits. These indices map directly to the appropriate rows and columns in the denoted matrix. Languages like Quipper [Green et al. 2013] and QWIRE [Paykin et al. 2017] instead use variables to refer to abstract qubits, which are later allocated to concrete ones. While aiding compositionality, this approach complicates formal proof. The semantics of a program that uses abstract qubits must convert those qubits into concrete indices. Reasoning about this conversion can be laborious, confounding automation, especially for recursive circuits. Moreover, notions like disjointness that are obvious when using concrete qubits—\(G_1 m\) operates on a different qubit than \(G_2 n\) when \(m \neq n\)—are no longer obvious when using abstract ones—\(G_1 x\) and \(G_2 y\) for variables \(x \neq y\) may not be disjoint if \(x\) and \(y\) could be allocated to the same concrete qubit. Disjointness is especially important in proofs of equivalence because gates acting on disjoint qubits commute, allowing us to reason about gates acting on different parts of the circuit in isolation. From a proof-engineering standpoint these benefits have been pivotal in allowing our proofs to scale up. Appendix A presents a detailed comparison of sqir and QWIRE, exploring the tradeoffs of concrete versus abstract qubits at a lower level.

**Example proof.** As an example of a proof we can carry out using sqir, we show that \(ghz\) , sqir’s Greenberger-Horne-Zeilinger (GHZ) state [Greenberger et al. 1989] preparation circuit given in Section 3.1.1, correctly produces the GHZ state. The GHZ state is an \(n\)-qubit entangled quantum state of the form \(\frac{1}{\sqrt{2}}(|0\rangle^{\otimes n} + |1\rangle^{\otimes n})\). This vector can be defined in Coq as shown at the right. Like our definition of \(|j\rangle^{\otimes n}\) discussed above, we declare that this expression has type \(\text{Vector} \ 2^n\), which will allow us to use it in any context where Coq expects a vector, deferring the proof that it is well-formed.

Our goal is to show that for any \(n > 0\) the circuit generated by \(ghz\ n\) produces the corresponding GHZ \(n\) vector when applied to \(|0\rangle^{\otimes n}\).

**Definition** GHZ \((n : \mathbb{N}) : \text{Vector} \ (2 \times n) :=\)

\[
\text{match } n \text{ with } \\
| 0 \Rightarrow I 1 \\
| s n' \Rightarrow \frac{1}{\sqrt{2}} \ast |0\rangle^{\otimes n} + \frac{1}{\sqrt{2}} \ast |1\rangle^{\otimes n}
\]

end.

**Lemma** \(ghz\text{\_correct} : \forall \ n : \mathbb{N},\)

\[n > 0 \rightarrow \parallel ghz \ n \parallel_n \times |0\rangle^{\otimes n} = \text{GHZ} \ n.\]

The proof proceeds by induction on \(n\). The \(n = 0\) case is trivial as it contradicts our hypothesis. For \(n = 1\) we show that \(H\) applied to \(|0\rangle\) produces the \(|+\rangle\) state. In the inductive step, the induction hypothesis says that the result of applying \(ghz\ n'\) to the input state \(nket \ n' \ 0\rangle\) is the state \((\frac{1}{\sqrt{2}} \ast |0\rangle^{\otimes n'} + \frac{1}{\sqrt{2}} \ast |1\rangle^{\otimes n'}) \otimes |0\rangle\). By applying \(\text{CNOT} \ (n' - 1) \ n'\) to this state, we show that \(ghz \ (n' + 1) = \text{GHZ} \ (n' + 1)\). Our use of concrete indices allows us to easily describe the semantics of \(\text{CNOT} \ x y\). If we had instead used abstract wires (e.g. variables \(x\) and \(y\)), then to reason about the semantics of \(\text{CNOT} \ x y\) we would also need to reason about the conversion of \(x\) and \(y\) to concrete indices, showing that in the inductive case \(x\) refers to a qubit in the GHZ state prepared by the recursive call and \(y\) references a fresh \(|0\rangle\) qubit.

Proof of this program is relatively simple, but proofs of properties about other programs can be more involved. Appendix B presents three additional programs: superdense coding with a proof
that it faithfully transmits the input bits, the Deutsch-Jozsa algorithm over $n$ qubits with a proof that it correctly distinguishes between a constant and balanced Boolean oracle, and quantum phase estimation over $n$ qubits with a proof that it correctly computes the phase of the input program.

### 3.2 Full sqIR: Adding Measurement

To describe general quantum programs $P$, we extend unitary sqIR with a *branching measurement* operation.

$$ P := \text{skip} \mid P_1 \mid P_2 \mid U \mid \text{meas } q P_1 P_2 $$

The command $\text{meas } q P_1 P_2$ (inspired by a similar construct in QPL [Selinger 2004]) measures the qubit $q$ and either performs program $P_1$ or $P_2$ depending on the result. We define non-branching measurement and resetting a qubit to $\mathbb{0} \equiv 0$ in terms of branching measurement:

- $\text{measure } q = \text{meas } q \text{ skip skip}$
- $\text{reset } q = \text{meas } q (X q) \text{ skip}$

Figure 4 defines the semantics of a non-unitary program as a function from density matrices $\rho$ to density matrices, following the approach of several previous efforts [Paykin et al. 2017; Ying 2011]. Density matrices provide a way to describe arbitrary quantum states, including mixed states which are probability distributions over quantum pure states and arise in the analysis of general quantum programs. For example, $\frac{1}{2} \left( \begin{array}{cc} 1 & 0 \\ 0 & 1 \end{array} \right)$ represents a 50% chance of $\mathbb{0}$ and a 50% chance of $\mathbb{1}$. (We also define a non-deterministic semantics in Appendix B, which is sometimes more convenient to use than density matrices.)

**Example: Quantum Teleportation.** The goal of quantum teleportation is to transmit a state $|\psi\rangle$ from one party (Alice) to another (Bob) using a shared entangled state. The circuit for quantum teleportation is shown in Figure 5 and the corresponding sqIR program is given below.

**Definition** bell : ucom base 3 := H 1 ; CNOT 1 2.
**Definition** alice : com base 3 := CNOT 0 1 ; H 0 ; measure 0 ; measure 1.
**Definition** bob : com base 3 := CNOT 1 2 ; CZ 0 2 ; reset 0 ; reset 1.
**Definition** teleport : com base 3 := bell ; alice ; bob.

The bell circuit prepares a Bell pair on qubits 1 and 2, which are respectively sent to Alice and Bob. Alice applies CNOT from qubit 0 to qubit 1 and then measures both qubits and (implicitly) sends them to Bob. Finally, Bob performs operations controlled by the (now classical) values on qubits 0 and 1 and then resets them to the zero state.

The correctness property for this program says that for any (well-formed) density matrix $\rho$, teleport takes the state $\rho \otimes |0\rangle\langle 0| \otimes |0\rangle\langle 0|$ to the state $|0\rangle\langle 0| \otimes |0\rangle\langle 0| \otimes \rho$.

**Lemma** teleport_correct : $\forall (\rho : \text{Density } 2)$, $WF_{\text{Matrix } \rho \rightarrow \{\text{teleport}\}_3 (\rho \otimes |0\rangle\langle 0| \otimes |0\rangle\langle 0|) = |0\rangle\langle 0| \otimes |0\rangle\langle 0| \otimes \rho$. 


Fig. 5. Circuit for quantum teleportation. In the standard presentation Bob only acts on the last qubit, given two classical bits as input. In our presentation, Bob (equivalently) performs operations controlled by the first two qubits, which are in a post-measurement classical state. We include the reset operations to simplify our statement of correctness.

The proof is simple: We perform (automated) arithmetic to show that the output matrix has the desired form.

Quantum teleportation is a rare case in which we prove something about a fixed-size program (i.e., \( n \) is fixed to 3 qubits) and not one of arbitrary dimension. Using the density matrix semantics to prove properties about programs of arbitrary dimension \( n \) is more involved. Stating the property requires introducing a symbolic density matrix \( \rho \), which will be multiplied on the left and right by \( 2^n \times 2^n \) matrices in the denotation. In our experience this results in complicated proof terms that are tedious to manipulate, even with automation. By contrast, when reasoning about the equivalence of two unitary programs we can simply compare their unitary matrices, without carrying around a symbolic \( \rho \) (or input vector). Such equivalences form the core of the \texttt{voqc} compiler, which leads us to favor the unitary semantics for our proofs.

4 OPTIMIZING UNITARY SQIR PROGRAMS

This section and the next describe \texttt{voqc}, our verified optimizer for quantum circuits. \texttt{voqc} primarily implements optimizations inspired by the state-of-the-art circuit optimizer of Nam et al. [2018]. As such, we do not claim credit for the optimizations themselves. Rather, our contribution is a framework that is sufficiently flexible that it can be used to prove such state-of-the-art optimizations correct. This section focuses on \texttt{voqc}'s optimizations on unitary SQIR programs; the next section discusses how \texttt{voqc} optimizes full (non-unitary) SQIR programs and maps them to connectivity-constrained architectures.

4.1 Program Equivalence

The \texttt{voqc} compiler takes as input a SQIR program and attempts to reduce its total gate count by applying a series of optimizations. \texttt{voqc}'s unitary optimizations are defined as Coq functions that map an input unitary SQIR program to an optimized one. For each optimization, we verify that it is semantics preserving (or sound), meaning that the output program is guaranteed to be equivalent to the input program.

We say that two unitary programs of dimension \( d \) are equivalent, written \( U_1 \equiv U_2 \), if their denotation is the same, i.e., \( \llbracket U_1 \rrbracket_d = \llbracket U_2 \rrbracket_d \). We also support a more general version of equivalence: We say that two circuits are equivalent up to a global phase, written \( U_1 \equiv U_2 \), when there exists a \( \theta \) such that \( \llbracket U_1 \rrbracket_d = e^{i\theta} \llbracket U_2 \rrbracket_d \). This is useful in the quantum setting because \( |\psi\rangle \) and \( e^{i\theta} |\psi\rangle \) (for \( \theta \in \mathbb{R} \)) represent the same physical state. Note that the latter notion of equivalence matches the former when \( \theta = 0 \).

Given this definition of equivalence we can write our soundness condition for optimization \texttt{opt} (which takes as input a program that uses an arbitrary gate set and number of qubits) as follows.

\texttt{Definition} sound \( \{ G \} \) (\texttt{opt} : \( \forall \{ d : \mathbb{N} \}, \text{ucom \ G d} \rightarrow \text{ucom \ G d} \) ) :=
\[ \forall (d : \mathbb{N}) (u : u_{\text{com}} G d), \llbracket \text{opt } u \rrbracket_d \cong \llbracket u \rrbracket_d. \]

This property is quantified over \( G, d, \) and \( u, \) meaning that the property holds for any program that uses any set of gates and any number of qubits. (The optimizations in our development are defined over a particular gate set, defined below, but still apply to programs that use any number of qubits. Our statements of soundness also occasionally have an additional precondition that requires program \( u \) to be well typed.)

### 4.2 \text{voqc} Optimization Overview

\text{voqc} implements two basic kinds of optimizations: replacement and propagate-cancel. The former simply identifies a pattern of gates and replaces it with an equivalent pattern. The latter works by commuting sets of gates when doing so produces an equivalent quantum program—often with the effect of “propagating” a particular gate rightward in the program—until two adjacent gates can be removed because they cancel out.

To ease the implementation of and proofs about these optimizations, we developed a framework of supporting library functions that operate on \text{sqir} programs as lists of gate applications, rather than on the native \text{sqir} representation. The conversion code takes a sequence of gate applications in the original \text{sqir} program and flattens it so that a program like \( \{ G_1 p; G_2 q \}; G_3 r \) is represented as the Coq list \([ G_1 p ; G_2 q ; G_3 r ]\). The denotation of the list representation is the denotation of its corresponding \text{sqir} program. Examples of the list operations our framework provides include:

- Finding the next gate acting on a qubit that satisfies some predicate \( f \).
- Propagating a gate using a set of cancellation and commutation rules (see Section 4.3).
- Replacing a sub-program with an equivalent program (see Section 4.4).
- Computing the maximal matching prefix of two programs.

We verify that these functions have the intended behavior (e.g., in the last example, that the returned sub-program is indeed a prefix of both input programs).

Our framework supports arbitrary gate sets. However, for \text{voqc} we chose to focus on a specific, universal gate set \( \{ H, X, R_z, \text{CNOT} \} \) where \( R_z(k) \) describes rotation about the z-axis by \( k \cdot \pi \) for \( k \in \mathbb{Q} \). Either the parser must produce input programs using this gate set, or \text{voqc} must convert the program to use it before optimizations can be applied. To compute a program’s denotation, \text{voqc}’s gates \( H, X, \) and \( R_z(k) \) are translated into \( R_{\pi/2, 0, \pi} , R_{\pi, 0, \pi} , \) and \( R_{0, 0, k\pi} \) in \text{sqir}’s base gate set. (\text{CNOT} translates to itself.)

Note that we chose not to parameterize rotations \( R_z \) by arbitrary reals because this would make verification unsound if these parameters were extracted to OCaml floating point numbers. Most existing tools (e.g. Qiskit [Aleksandrowicz et al. 2019] and Nam et al. [2018]) allow gates parameterized by floats, which invites rounding error and can lead to unsound optimization. \text{voqc}’s gate set is identical to Nam et al.’s, with the exception of the z-axis rotation parameter type.

### 4.3 Optimization by Propagation and Cancellation

Our propagate-cancel optimizations have two steps. First we localize a set of gates by repeatedly applying commutation rules. Then we apply a circuit equivalence to replace that set of gates. In \text{voqc} most optimizations of this form use a library of code patterns, but one—not propagation—is slightly different, so we discuss it first.

**Not Propagation.** The goal of not propagation is to remove cancelling \( X \) (“not”) gates. Two \( X \) gates cancel when they are adjacent or they are separated by a circuit that commutes with \( X \). We find \( X \) gates separated by commuting circuits by repeatedly applying the propagation rules in Figure 6. An example application of the not propagation algorithm is shown in Figure 7.
A Verified Optimizer for Quantum Circuits

\[
\begin{align*}
X q; H q & \equiv H q; Z q \\
X q; Rz(k) q & \equiv Rz(2-k) q; X q \\
X q_1; CNOT q_1 q_2 & \equiv CNOT q_1 q_2; X q_1; X q_2 \\
X q_2; CNOT q_1 q_2 & \equiv CNOT q_1 q_2; X q_2 \\
\end{align*}
\]

Fig. 6. Equivalences used in not propagation.

Fig. 7. An example of not propagation. In the first step the leftmost \(X\) gate propagates through the \(CNOT\) gate to become two \(X\) gates. In the second step the upper \(X\) gate propagates through the \(H\) gate and the lower \(X\) gates cancel.

This implementation may introduce extra \(X\) gates at the end of a circuit or extra \(Z\) gates in the interior of the circuit. Extra \(Z\) gates are likely to be cancelled by the gate cancellation and rotation merging passes that follow, and moving \(X\) gates to the end of a circuit makes the rotation merging optimization more likely to succeed.

We note that our version of this optimization is a simplification of Nam et al.’s, which is specialized to a three-qubit \(TOFF\) gate; this gate can be decomposed into a \(\{H, Rz, CNOT\}\) program. In our experiments, we did not observe any difference in performance between \(voqc\) and Nam et al. due to this simplification.

**Gate Cancellation.** The single- and two-qubit gate cancellation optimizations rely on the same propagate-cancel pattern used in not propagation, except that gates are returned to their original location if they fail to cancel. To support this pattern, we provide a general \(\text{propagate}\) function in voqc. This function takes as inputs (i) an instruction list, (ii) a gate to propagate, and (iii) a set of rules for commuting and cancelling that gate. At each iteration, \(\text{propagate}\) performs the following actions:

1. Check if a cancellation rule applies. If so, apply that rule and return the modified list.
2. Check if a commutation rule applies. If so, commute the gate and recursively call \(\text{propagate}\) on the remainder of the list.
3. Otherwise, return the gate to its original position.

We have proved that our \(\text{propagate}\) function is sound when provided with valid commutation and cancellation rules.

Each commutation or cancellation rule is implemented as a partial Coq function from an input circuit to an output circuit. A common pattern in these rules is to identify one gate (e.g., an \(X\) gate), and then to look for an adjacent gate it might commute with (e.g., \(CNOT\)) or cancel with (e.g., \(X\)). For commutation rules, we use the rewrite rules shown Figure 8. For cancellation rules, we use the fact that \(H, X,\) and \(CNOT\) are all self-cancelling and \(Rz(k)\) and \(Rz(k')\) combine to become \(Rz(k + k')\).

### 4.4 Circuit Replacement

We have implemented two optimizations—Hadamard reduction and rotation merging—that work by replacing one pattern of gates with an equivalent one; no preliminary propagation is necessary. These aim either to reduce the gate count directly, or to set the stage for additional optimizations.
Fig. 8. Commutation equivalences for single- and two-qubit gates adapted from Nam et al. [2018, Figure 5]. We use the second and third rules for propagating both single- and two-qubit gates.

Fig. 9. Equivalences for removing Hadamard gates adapted from Nam et al. [2018, Figure 4]. $P$ is the phase gate and $P^\dagger$ is its inverse.

**Hadamard Reduction.** The Hadamard reduction routine employs the equivalences shown in Figure 9 to reduce the number of $H$ gates in the program. Removing $H$ gates is useful because $H$ gates limit the size of the $\{Rz, CNOT\}$ subcircuits used in the rotation merging optimization.

**Rotation Merging.** The rotation merging optimization allows for combining $Rz$ gates that are not physically adjacent in the circuit. This optimization is more sophisticated than the previous optimizations because it does not rely on small structural patterns (e.g., that adjacent $X$ gates cancel), but rather on more general (and non-local) circuit behavior. The basic idea behind rotation merging is to (1) identify subcircuits consisting of only $CNOT$ and $Rz$ gates and (2) merge $Rz$ gates within those subcircuits that are applied to qubits in the same logical state.

The argument for the correctness of this optimization relies on the *phase polynomial* representation of a circuit. Let $C$ be a circuit consisting of $CNOT$ gates and rotations about the $z$-axis. Then on basis state $|x_1, ..., x_n\rangle$, $C$ will produce the state

$$e^{i p(x_1, ..., x_n)} |h(x_1, ..., x_n)\rangle$$
where $h : \{0, 1\}^n \to \{0, 1\}^n$ is an affine reversible function and

$$p(x_1, \ldots, x_n) = \sum_{i=1}^l (\theta_i \mod 2\pi) f_i(x_1, \ldots, x_n)$$

is a linear combination of affine boolean functions. $p(x_1, \ldots, x_n)$ is called the phase polynomial of circuit $C$. Each rotation gate in the circuit is associated with one term of the sum and if two terms of the phase polynomial satisfy $f_i(x_1, \ldots, x_n) = f_j(x_1, \ldots, x_n)$ for some $i \neq j$, then the corresponding $i$ and $j$ rotations can be merged.

As an example, consider the two circuits shown below.

To prove that these circuits are equivalent, we can consider their behavior on basis state $|x_1, x_2\rangle$. Recall that applying $Rz(k)$ to the basis state $|x\rangle$ produces the state $e^{ik\pi x} |x\rangle$ and CNOT $|x, y\rangle$ produces the state $|x, x \oplus y\rangle$ where $\oplus$ is the xor operation. Evaluation of the left-hand circuit proceeds as follows:

$$|x_1, x_2\rangle \to e^{ik\pi x_2} |x_1, x_2\rangle$$
$$\quad \to e^{ik\pi x_2} |x_1, x_1 \oplus x_2\rangle$$
$$\quad \to e^{ik\pi x_2} |x_2, x_1 \oplus x_2\rangle$$
$$\quad \to e^{ik\pi x_2} e^{ik'\pi x_2} |x_2, x_1 \oplus x_2\rangle.$$  

Whereas evaluation of the right-hand circuit produces

$$|x_1, x_2\rangle \to |x_1, x_1 \oplus x_2\rangle$$
$$\quad \to |x_2, x_1 \oplus x_2\rangle$$
$$\quad \to e^{i(k+k')\pi x_2} |x_2, x_1 \oplus x_2\rangle.$$  

The two resulting states are equal because $e^{ik\pi x_2} e^{ik'\pi x_2} = e^{i(k+k')\pi x_2}$. This implies that the unitary matrices corresponding to the two circuits are the same. We can therefore replace the circuit on the left with the one on the right, removing one gate from the circuit.

Our rotation merging optimization follows the reasoning above for arbitrary $\{Rz, CNOT\}$ circuits. For every gate in the program, it tracks the Boolean function associated with every qubit (the Boolean functions above are $x_1, x_2, x_1 \oplus x_2$), and merges $Rz$ rotations when they are applied to qubits associated with the same Boolean function. To prove equivalence over $\{Rz, CNOT\}$ circuits, we show that the original and optimized circuits produce the same output on every basis state. We have found evaluating behavior on basis states is useful for proving equivalences that are not as direct as those listed in Figures 8 and 9.

Although our merge operation is identical to Nam et al.’s, our approach to constructing $\{Rz, CNOT\}$ subcircuits differs. We construct a $\{Rz, CNOT\}$ subcircuit beginning from a $Rz$ gate whereas Nam et al. begin from a $CNOT$ gate. The result of this simplification is that we may miss some opportunities for merging. However, in our experiments (Section 6) we found that this choice impacted only one benchmark.

### 4.5 Proving Low-Level Circuit Equivalences

VOQC optimizations make heavy use of circuit equivalences such as those shown in Figures 6, 8 and 9. So to prove that VOQC optimizations are sound, we must formally verify these equivalences are correct. Such proofs require showing equality between two matrix expressions, which can be
tedious in the case where the matrix size is left symbolic. For example, consider the following equivalence used in not propagation:

\[ X \; n; \; \text{CNOT} \; m \; n \equiv \text{CNOT} \; m \; n; \; X \; n \]

for arbitrary \( n, m \) and dimension \( d \). Applying our definition of equivalence, this amounts to proving

\[
apply_1(X, n, d) \times apply_2(\text{CNOT}, m, n, d) = apply_2(\text{CNOT}, m, n, d) \times apply_1(X, n, d),
\]

per Figure 3. Suppose both sides of the equation are well typed (\( m < d \) and \( n < d \) and \( m \neq n \)), and consider the case where \( m < n \) (the \( n < m \) case is similar). We expand \( apply_1 \) and \( apply_2 \) as follows with \( p = n - m - 1 \) and \( q = d - n - 1 \):

\[
apply_1(X, n, d) = \mathbb{I}_m \otimes \sigma_x \otimes \mathbb{I}_q
\]

\[
apply_2(\text{CNOT}, m, n, d) = \mathbb{I}_m \otimes |1\rangle \langle 1| \otimes \mathbb{I}_p \otimes \sigma_x \otimes \mathbb{I}_q + \mathbb{I}_m \otimes |0\rangle \langle 0| \otimes \mathbb{I}_p \otimes \mathbb{I}_q
\]

Here, \( \sigma_x \) is the matrix interpretation of the \( X \) gate and \(|1\rangle \langle 1| \otimes \sigma_x + |0\rangle \langle 0| \otimes \mathbb{I}_q \) is the matrix interpretation of the \( \text{CNOT} \) gate (in Dirac notation). We complete the proof of equivalence by normalizing and simplifying each side of Equation (1), showing both sides to be the same.

Automation. We address the tedium of such \( \text{voqc} \) proofs by almost entirely automating the matrix normalization and simplification steps. We provide a Coq tactic called \texttt{gridify} for proving general equivalences correct. Rather than assuming \( m < n < d \) as above, the \texttt{gridify} tactic does case analysis, immediately solving all cases where the circuit is ill-typed (e.g., \( m = n \) or \( d < m \)) and thus has the zero matrix as its denotation. In the remaining cases \( (m < n \) and \( n < m \) above), it puts the expressions into a form we call \textit{grid normal} and applies a set of matrix identities.

In grid normal form, each arithmetic expression has addition on the outside, followed by tensor product, with multiplication on the inside, i.e., \( ((\ldots \otimes (\ldots \otimes (\ldots \otimes (\ldots \otimes (\ldots \otimes (\ldots \otimes (\ldots \otimes (\ldots \otimes (\ldots \otimes \ldots \otimes \ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots )\ldots ). \) The \texttt{gridify} tactic rewrites an expression into this form by using the following rules of matrix arithmetic:

- \( I_{mn} = I_m \otimes I_n \)
- \( A \times (B + C) = A \times B + A \times C \)
- \( (A + B) \times C = A \times C + B \times C \)
- \( A \otimes (B + C) = A \otimes B + A \otimes C \)
- \( (A + B) \otimes C = A \otimes C + B \otimes C \)
- \( (A \otimes B) \times (C \otimes D) = (A \times C) \otimes (B \times D) \)

The first rule is applied to facilitate application of the other rules. (For instance, in the example above, \( \mathbb{I}_m \) would be replaced by \( \mathbb{I}_m \otimes \mathbb{I}_2 \otimes \mathbb{I}_p \) to match the structure of the \( \text{apply}_2 \) term.) After expressions are in grid normal form, \texttt{gridify} simplifies them by removing multiplication by the identity matrix and rewriting simple matrix products (e.g. \( \sigma_x \sigma_x = \mathbb{I}_2 \)).

In our example, normalization and simplification by \texttt{gridify} rewrites each side of the equality in Equation (1) to be the following

\[
\mathbb{I}_m \otimes |1\rangle \langle 1| \otimes \mathbb{I}_p \otimes \mathbb{I}_q + \mathbb{I}_m \otimes |0\rangle \langle 0| \otimes \mathbb{I}_p \otimes \sigma_x \otimes \mathbb{I}_q,
\]

thus proving that the two expressions are equal.

We use \texttt{gridify} to verify most of the equivalences used in the optimizations given in Sections 4.3 and 4.4. The tactic is most effective when equivalences are small: The equivalences used in \textit{gate cancellation} and \textit{Hadamard reduction} apply to patterns of at most five gates applied to up to three qubits within an arbitrary circuit. For equivalences over large sets of qubits, like the one used in \textit{rotation merging}, we do not use \texttt{gridify} directly, but still rely on our automation for matrix simplification.
4.6 Scheduling
The \texttt{voqc optimize} function applies each of the optimizations we have discussed one after the other, in the following order (due to Nam et al.):

$$0, 1, 3, 2, 3, 1, 2, 4, 3, 2$$

where 0 is not propagation, 1 is Hadamard reduction, 2 is single-qubit gate cancellation, 3 is two-qubit gate cancellation, and 4 is rotation merging. The rationale for this ordering is that removing $X$ and $H$ gates (0,1) allows for more effective application of the gate cancellation (2,3) and rotation merging (4) optimizations. In our experiments (Section 6), we observed that single-qubit gate cancellation and rotation merging were the most effective at reducing gate count.

5 OTHER VERIFIED TRANSFORMATIONS
We have also implemented verified optimizations of non-unitary programs in \texttt{voqc} (inspired by optimizations in IBM’s Qiskit compiler [Aleksandrowicz et al. 2019]) and verified a transformation that maps a circuit to a connectivity-constrained architecture.

5.1 Non-unitary Optimizations
We have implemented two non-unitary optimizations: removing pre-measurement $z$-rotations, and classical state propagation. For these optimizations, we represent a non-unitary program $P$ as a list of \textit{blocks}. A block is a binary tree whose leaves are unitary programs (in list form) and nodes are measurements $\texttt{meas} \; q \; P_1 \; P_2$ whose children $P_1$ and $P_2$ are lists of blocks. Since the density matrix semantics denotes programs as functions over matrices, we say that programs $P_1$ and $P_2$ of dimension $d$ are equivalent if for every input $\rho$, $\mathcal{P}_1 \mathcal{P}_2 \mathcal{P}_0 (\rho) = \mathcal{P}_1 \mathcal{P}_2 \mathcal{P}_0 (\rho)$.

$z$-rotations Before Measurement. $z$-axis rotations (or, more generally, diagonal unitary operations) before a measurement will have no effect on the measurement outcome, so they can safely be removed from the program. The optimization locates $Rz$ gates before measurement operations and removes them. This optimization was inspired by the RemoveDiagonalGatesBeforeMeasure pass implemented in Qiskit.

Classical State Propagation. Once a qubit has been measured, the subsequent branch taken provides information about the qubit’s (now classical) state, which may allow pre-computation of some values. For example, in the branch where qubit $q$ has been measured to be in the $|0\rangle$ state, any \texttt{CNOT} with $q$ as the control will be a no-op and any subsequent measurements of $q$ will still produce zero.

In detail, given a qubit $q$ in classical state $|i\rangle$, our optimization applies these propagation rules:

- $Rz(k) \; q$ preserves the classical state of $q$.
- $X \; q$ flips the classical state of $q$.
- If $i = 0$ then \texttt{CNOT} $q \; q'$ is removed, and if $i = 1$ then \texttt{CNOT} $q \; q'$ becomes $X \; q'$.
- $\texttt{meas} \; q \; P_1 \; P_0$ becomes $P_1$.
- $H \; q$ and \texttt{CNOT} $q' \; q$ make $q$ non-classical and terminate analysis.

Our statement of correctness for one round of propagation says that if qubit $q$ is in a classical state in the input, then the optimized program will have the same denotation as the unoptimized original. We express the requirement that qubit $q$ be in classical state $i \in \{0, 1\}$ with the condition $|i\rangle_q \langle i| \times \rho \times |i\rangle_q \langle i| = \rho$, which says that projecting state $\rho$ onto the subspace where $q$ is in state $|i\rangle$ results in no loss of information.
This optimization is not implemented directly in Qiskit, but Qiskit contains passes that have a similar effect. For example, the RemoveResetInZeroState pass removes adjacent reset gates, as the second has no effect.

5.2 Circuit Mapping

Similar to how optimization aims to reduce qubit and gate usage to make programs more feasible to run on near-term machines, circuit mapping aims to address the connectivity constraints of near-term machines [Saeedi et al. 2011; Zulehner et al. 2017]. Circuit mapping algorithms take as input an arbitrary circuit and output a circuit that respects the connectivity constraints of some underlying architecture.

For example, consider the connectivity of IBM’s five-qubit Tenerife machine shown in Figure 10. This is a representative example of a superconducting qubit system, where qubits are laid out in a 2-dimensional grid and possible interactions are described by directed edges between the qubits. The direction of the edge indicates which qubit can be the control of a two-qubit gate and which can be the target. For instance, a CNOT gate may be applied with Q4 as the control and Q2 as the target, but not the reverse. However, no two-qubit gate is possible between physical qubits Q4 and Q1 on the Tenerife.

We have implemented a simple circuit mapper for unitary $\text{sqir}$ programs and verified that it is sound and produces programs that satisfy the relevant hardware constraints. Our circuit mapper is parameterized by two functions that describe the connectivity of an architecture: one function determines whether an edge is in the connectivity graph and another function finds an undirected path between any two nodes. Our mapping algorithm takes as input these functions, a program referencing logical qubits, and a map expressing the initial correspondence between the program’s logical qubits and the physical qubits available on the architecture. The algorithm produces a program referencing physical qubits as well as an updated correspondence. Every time a CNOT occurs between two logical qubits whose corresponding physical qubits are not adjacent in the underlying architecture, we insert SWAP operations to move the target and control into adjacent positions and update the physical-logical qubit correspondence accordingly. To apply a CNOT when an edge points in the wrong direction, we make use of the equivalence $\text{CNOT}_b a \equiv \text{H} a \text{H} b; \text{CNOT}_a b; \text{H} a; \text{H} b$. For soundness, we prove that the mapped circuit is equivalent to the original up to a permutation of the qubits.

Although our mapping algorithm is simple, it allows for some flexibility in design because we do not specify the method for choosing the initial physical-logical qubit correspondence (called “placement”) or the implementation of the function that finds paths in the connectivity graph (“routing”). This allows, for example, placement and routing strategies that take into account error characteristics of the machine [Tannu and Qureshi 2019]. We expect that our verification framework can be applied to more sophisticated mapping algorithms such as those that partition the circuit into layers and insert SWAPs between layers rather than naively inserting SWAPs before CNOT gates [Zulehner et al. 2017]. We have implemented and verified mapping functions for the Tenerife
6 EXPERIMENTAL EVALUATION

The value of voqc (and sqir) is determined by the quality of the verified optimizations we can write with it. We can judge optimization quality empirically. In particular, we can run voqc on a benchmark of circuit programs and see how well it optimizes those programs, compared to (non-verified) state-of-the-art compilers.

To this end, we compared the performance of voqc’s verified (unitary) optimizations against IBM’s Qiskit transpiler [Aleksandrowicz et al. 2019], CQC’s tket compiler [Cambridge Quantum Computing Ltd 2019; Sivarajah et al. 2020], PyZX [Kissinger and van de Wetering 2019], and the optimizers presented in Nam et al. [2018] and Amy et al. [2013] on a set of benchmarks developed by Amy et al. We find that voqc has comparable performance to all of these: it generally beats all but Nam et al. in terms of both total gate count reduction and T-gate reduction, and often matches Nam et al. However, our aim is not to claim superiority over these tools (after all, we have implemented a subset of the optimizations available in Nam et al., and Qiskit and tket contain many features that Nam et al. does not), but to demonstrate that the optimizations we have implemented are on par with existing unverified tools.

We also evaluated the performance of voqc on the full set of benchmarks used by Nam et al., of which Amy et al. is a subset. The results reported in this section are representative of voqc’s behavior on the full set, which is reported in Appendix C.

Benchmarks. We evaluated performance by applying voqc and the other compilers to the benchmark of Amy et al. [2013], which consists of programs written in the “Clifford+T” gate set. We measure the reduction in total gate count and T-gate count. Total gate count is a useful metric for near-term quantum computing, where the length of the computation must be minimized to reduce error. T-gate count is relevant in the fault-tolerant regime where qubits are encoded using quantum error correcting codes and operations are performed fault-tolerantly. In this regime the standard method for making Clifford+T circuits fault tolerant produces particularly expensive translations for T gates, so reducing T-count is a common optimization goal. The Clifford+T set is a subset of voqc’s gate set where each z-axis rotation is restricted to be a multiple of \( \pi/4 \) (an odd multiple of \( \pi/4 \) corresponds to one T gate).

The benchmarks consist of arithmetic circuits and implementations of multiple-control Toffoli gates, ranging from 45 to 13,593 gates and 5 to 96 qubits. These benchmarks only contain unitary circuits, so they only serve to evaluate our unitary circuit optimizations. We do not evaluate circuit mapping transformations. Some of the benchmarks contain the three-qubit CCZ gate. Before applying optimizations we convert CCZ gates to voqc’s gate set using the following standard decomposition, where T is the Rz(1/4) gate and T\( ^\dagger \) is its inverse Rz(7/4).

\[
\text{CCZ a b c := [ CNOT b c ; T}^\dagger \text{ c ; CNOT a c ; T c ; CNOT b c ; T}^\dagger \text{ c ; CNOT a c ; CNOT a b ; T}^\dagger \text{ b ; CNOT a b ; T a ; T b ; T c ]}.
\]
Table 1. Summary of unitary optimizations for reducing total gate count.

| Optimization Type                                      | Qiskit 0.17.0                          | t|ket) 0.4.2                          | voqc
|-------------------------------------------------------|---------------------------------------|---------------------------------|--------
| Not propagation (P)                                   | ✓                                     | ✓                               | ✓      |
| Hadamard gate reduction (L, H)                        | ✓                                     | ✓                               | ✓      |
| Single-qubit gate cancellation (L, H)                 | ✓                                     | ✓                               | ✓      |
| Two-qubit gate cancellation (L, H)                    | ✓                                     | ✓                               | ✓      |
| Rotation merging using phase polynomials (L)          | ✓                                     | ✓                               | ✓      |
| Floating $R_z$ gates (H)                               |                                       |                                 |        |
| Special-purpose optimizations (L, H)                  |                                       |                                 | ✓      |
| • LCR optimizer                                       |                                       |                                 |        |
| • Toffoli decomposition                               |                                       |                                 |        |

Baseline: Total Gate Count. To evaluate reduction in total gate count, we compare voqc’s performance with that of Nam et al., Qiskit version 0.17.0, and t|ket) version 0.4.2. We do not include the results from Amy et al. or PyZX because their optimizations are aimed at reducing $T$-count, and often result in a higher total gate count. Table 1 performs a direct comparison of functionality provided by voqc versus Nam et al., Qiskit, and t|ket). For the Qiskit optimizations, $L_i$ indicates that a routine is used by optimization level $i$. For Nam et al., P stands for “preprocessing” and L and H indicate whether the routine is in the “light” or “heavy” versions of the optimizer. voqc provides the complete and verified functionality of the routines marked with ✓; we write ✓* to indicate that voqc contains a verified optimization with similar, although not identical, behavior.

Compared to Nam et al.’s rotation merging, voqc performs a slightly less powerful optimization (as discussed in Section 4.4). Conversely, voqc’s one- and two-qubit gate cancellation routines generalize Qiskit’s Optimize1qGates and CXCancellation when restricted to voqc’s gate set. For CommutativeCancellation, Qiskit’s routine follows the same pattern as our gate cancellation routines, but uses matrix multiplication to determine whether gates commute while we use a rule-based approach; neither is strictly more effective than the other. t|ket)’s FullPeepholeOptimise performs local rewrites similar to those applied by Qiskit.

When evaluating Qiskit, we include all unitary optimizations up to level 3. In our evaluation, both Qiskit and t|ket) use the gate set $\{u_1, u_2, u_3, \text{CNOT}\}$ where $u_3$ is $R_{\theta, \phi, \lambda}$ from voqc’s base set and $u_1$ and $u_2$ are $u_3$ with certain arguments fixed. This gate set gives these industrial compilers an advantage over voqc because they can, for example, represent $X$ followed by $H$ with a single gate.

Baseline: $T$-Gate Count. To evaluate reduction in $T$-gate count, we compare voqc against Nam et al., Amy et al., and PyZX. We do not include results from Qiskit or t|ket) because these compilers produce circuits that do not use the Clifford+T set. When evaluating PyZX, we use the simplify.full_reduce method, which applies an optimization similar in intent to rotation merging, but implemented in terms of the ZX-calculus. We do not re-run Nam et al. (which is proprietary software) or Amy et al., but rather report results from Nam et al. [2018]. We only report results of Nam et al.’s heavy optimization.
<table>
<thead>
<tr>
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<th>ket</th>
<th>Nam</th>
<th>VOQC</th>
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</tbody>
</table>

**Average Reduction** | – | 10.4% | 10.9% | 26.4% | 18.7% |

Table 2. Reduced total gate counts on Amy et al. [2013] benchmarks. Red cells indicate programs optimized incorrectly.

**Results.** The results are shown in Table 2 and Table 3. In each row, we have marked in bold the gate count of the best-performing optimizer. The average reduction for each optimizer is given in the last row, although performance varies substantially between benchmarks. Shaded cells mark that the resulting optimized circuit has been found to be inequivalent to the original circuit [Kissinger and van de Wetering 2019; Nam et al. 2018], indicating a bug in the relevant optimizer. Results for incorrectly-optimized circuits are not included in the averages on the last line.

On average, Qiskit reduces the total gate count by 10.4%, Q|ket| by 10.9%, Nam et al. by 26.4%, and VOQC by 18.7%. VOQC outperforms or matches the performance of Qiskit and Q|ket| on all benchmarks but one. In 8 out of 28 cases VOQC outperforms Nam et al. The gap in performance between VOQC and the industrial compilers is due to VOQC’s rotation merging optimization, which has no analogue in Qiskit or Q|ket|. The gap in performance between Nam et al. and VOQC is due to the fact that we have not yet implemented all their optimization passes (per Table 1). In particular,
<table>
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<th>PyZX</th>
<th>Nam</th>
<th>VOQC</th>
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</table>

Table 3. Reduced $T$-gate counts on the Amy et al. [2013] benchmarks. Red cells indicate programs optimized incorrectly.

Nam et al.’s “special-purpose Toffoli decomposition” (which affects how CCZ gates are decomposed) enables rotation merging and single-qubit gate cancellation to cancel two gates (e.g. cancel $T$ and $T^\dagger$) where we instead combine two gates into one (e.g. $T$ and $T$ becomes $P$). Interestingly, the cases where voqc outperforms Nam et al. can also be attributed to their Toffoli decomposition heuristics, which sometimes result in fewer cancellations than the naïve decomposition that we use. We expect that more sophisticated Toffoli decomposition can be verified in Coq, although at present it is difficult to reason about the correctness of Toffoli and CCZ circuits using our matrix library.

voqc’s performance is closer to Nam et al.’s when considering $T$-count. On average, Amy et al. reduce the $T$-gate count by 40.9%, PyZX by 43.8%, and Nam et al. and voqc by 42.3%. voqc matches Nam et al. on all benchmarks but two. The first case (qcla_adder_10) is due to our simplification in rotation merging. In the second case (qcla_mod_7), Nam et al.’s optimized circuit was later found to be inequivalent to the original circuit [Kissinger and van de Wetering 2019], so the lower $T$-count is
spurious. On 16 benchmarks, all optimizers produce the same $T$-count. This is somewhat surprising since, although all these optimizers rely on some form of rotation merging, their implementations differ substantially. Kissinger and van de Wetering [2019] posit that these results indicate a local optimum in the ancilla-free case for some of the benchmarks (in particular the tof benchmarks, whose $T$-count is not reduced by applying additional techniques [Heyfron and T. Campbell 2018]).

To compare the runtimes of the different tools, we ran voqc, Qiskit, tket, and PyZX on a standard laptop with a 2.9 GHz Intel Core i5 processor and 16 GB of 1867 MHz DDR3 memory, running macOS Mojave. We consider the timings for Amy et al. and Nam et al. given in Nam et al. [2018, Table 4], which were measured on a similar machine with 8 GB RAM running OS X El Capitan. We show the geometric mean runtimes over all 28 benchmarks below.

<table>
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<th>Tool</th>
<th>Nam (L)</th>
<th>Nam (H)</th>
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<th>tket</th>
<th>Amy</th>
<th>PyZX</th>
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<tbody>
<tr>
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<td>0.012s</td>
<td>0.002s</td>
<td>0.018s</td>
<td>2.128s</td>
<td>0.226s</td>
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</table>

All the tools are fast; Qiskit is the slowest at around 2 seconds. However, these means are not the entire story; the tools’ performances scale differently with increasing qubit and gate count. For example, on gf2^32_mult (the largest benchmark) Qiskit is faster than Nam et al. heavy optimization, tket, and Amy et al. For more details on voqc’s performance, see Appendix C.

These results are encouraging evidence that voqc supports useful and interesting verified optimizations, and that we have faithfully implemented Nam et al.’s optimizations. Furthermore, despite having been written with verification in mind, voqc’s runtime performance is not significantly worse than (and sometimes better than) that of current tools.

**Trusted Code.** For performance, voqc uses OCaml primitives for describing rational numbers, maps and sets, rather than the code extracted from Coq. Thus we implicitly trust that the OCaml implementation of these data types is consistent with Coq’s; we believe that this is a reasonable assumption. Furthermore, our translation from OpenQASM to sqir and extraction from Coq to OCaml are unverified.

7 RELATED WORK

**Verified Quantum Programming.** We designed sqir primarily as the intermediate language for voqc’s verified optimizations, but we find it adequate for verified source programming as well (per Section 3 and Appendix B).

Several lines of work have explored formally verifying aspects of a quantum computation. The earliest attempts to do so in a proof assistant were an Agda implementation of the Quantum IO Monad [Green 2010] and a small Coq quantum library by Boender et al. [2015]. These were both proofs of concept, and neither developed beyond verifying basic protocols.

The higher-level Qwire programming language is, like sqir, embedded in the Coq proof assistant, and has been used to verify a variety of simple programs [Rand et al. 2017], assertions regarding ancilla qubits [Rand et al. 2018b], and its own metatheory [Rand 2018]. voqc and sqir reuse parts of Qwire’s Coq development, and take inspiration and lessons from its design. However, as discussed in Section 3.1.3 and Appendix A, Qwire’s higher-level abstractions complicate verification. Moreover, such abstractions do not reflect the kind of quantum programming we can expect to do in the near future. For example, a key element of Qwire is dynamic lifting, which permits measuring a qubit and using the result as a Boolean value in the host language to compute the remainder of a circuit [Green et al. 2013]. Today’s quantum computers cannot reliably exchange information between a (typically supercooled) quantum chip and a classical computer before qubits decohere. Thus, practically-minded languages like IBM’s OpenQASM [Cross et al. 2017] only allow for a limited form of branching that is close to sqir’s.
Concurrently with this work, Chareton et al. [2020] introduced Qbricks, a tool implemented in Why3 [Filliâtre and Paskevich 2013] whose aim is to support mostly-automated verification of complex quantum algorithms. Their design in many ways mirrors sQIR’s: both tools provide special support for reasoning about quantum programs and the languages are simplified so that programs have a straightforward translation to their semantics. Qbricks primarily uses a path-sum semantics [Amy 2018] and only supports unitary programs. sQIR programs have a natural translation to their matrix semantics because sQIR uses natural numbers in place of variables. Qbricks removes variables from the language entirely, instead constructing circuits as objects via parallel and sequential composition. This, together with sQIR’s support for branching measurement, puts sQIR closer to languages like OpenQASM, which makes it well-suited to be an intermediate representation in a quantum compiler stack. A key focus of Qbricks is automation; we hope to provide better automation and more tactic support for voqc in the future, but we can expect that Coq proofs will require more manual effort than proofs in Why3. That said, sQIR’s design is suited to general verification; for example, we were able to verify quantum phase estimation, a key achievement of Qbricks, using sQIR (Appendix B).

Another line of work, pioneered by D’Hondt and Panangaden [2006] and Ying [2011], uses program logics to reason about quantum programs. These logics allow proof of a variety of program properties inside a formal deductive system. Liu et al. [2019] implemented Ying’s quantum Hoare logic inside the Isabelle proof assistant and used it to prove the correctness of Grover’s algorithm, and Unruh [2019] developed a relational quantum Hoare logic and built an Isabelle-based tool to prove the security of quantum cryptosystems. Implementing these kinds of logics in Coq and proving them correct with respect to sQIR’s denotational semantics may prove useful (though we have verified several interesting sQIR programs correct directly).

Verified Quantum Compilation. Quantum compilation is an active area. In addition to Qiskit, |ket⟩, and Nam et al. [2018] (discussed in Section 6), other recent compiler efforts include quilc [Rigetti Computing 2019b], ScaffCC [Javadi-Abhari et al. 2014], and Project Q [Steiger et al. 2018]. Due to resource limits on near-term quantum computers, most compilers for quantum programs contain some degree of optimization, and nearly all place an emphasis on satisfying architectural requirements, like mapping to a particular gate set or qubit topology. None of the optimization or mapping code in these compilers is formally verified.

However, voqc is not the only quantum compiler to which automated reasoning or formal verification has been applied. Amy et al. [2017] developed a certified optimizing compiler from source Boolean expressions to reversible circuits, but did not handle general quantum programs. Rand et al. [2018b] developed a similar compiler for quantum circuits but without optimizations (using the Qwire language).

The problem of optimization verification has also been considered in the context of the ZX-calculus [Coecke and Duncan 2011], which is a formalism for describing quantum tensor networks (which generalize quantum circuits) based on categorical quantum mechanics [Abramsky and Coecke 2009]. The ZX-calculus is characterized by a small set of rewrite rules that allow translation of a diagram to any other diagram representing the same computation [Jeandel et al. 2018]. Fagan and Duncan [2018] verified an optimizer for ZX diagrams representing Clifford circuits (which use the non-universal gate set \{CNOT, H, S\}) in the Quantomatic graphical proof assistant [Kissinger and Zamdzhiev 2015]. PyZX [Kissinger and van de Wetering 2019] uses ZX diagrams as an intermediate representation for compiling quantum circuits, and generally achieves performance comparable to leading compilers [Kissinger and van de Wetering 2019]. While PyZX is not verified in a proof assistant like Coq (the “Py” stands for Python), it does rely on a small, well-studied equational theory. Additionally, PyZX can perform translation validation to check if a compiled circuit is
equivalent to the original. However, PyZX’s translation validator is not guaranteed to succeed for any two equivalent circuits.

A recent paper from Smith and Thornton [2019] presents a compiler with built-in translation validation via QMDD equivalence checking [Miller and Thornton 2006]. However the optimizations they consider are much simpler than voqc’s and the QMDD approach scales poorly with increasing number of qubits. Our optimizations are all verified for arbitrary dimension.

Concurrently with our work, Shi et al. [2019] developed CertiQ, an approach to verifying properties of circuit transformations in the Qiskit compiler, which is implemented in Python. Their approach has two steps. First, it uses matrix multiplication to check that the unitary semantics of two concrete gate patterns are equivalent. Second, it uses symbolic execution to generate verification conditions for parts of Qiskit that manipulate circuits. These are given to an SMT solver to verify that pattern equivalences are applied correctly according to programmer-provided function specifications and invariants. That CertiQ can analyze Python code directly in a mostly automated fashion is appealing. However, it is limited in the optimizations it can verify. For example, equivalences that range over arbitrary indices, like $\text{CNOT } m \ x; \ \text{CNOT } n \ x \equiv \text{CNOT } n \ x; \ \text{CNOT } m \ x$ cannot be verified by matrix multiplication; CertiQ checks a concrete instance of this pattern and then applies it to more general circuits. More complex optimizations like rotation merging (the most powerful optimization in our experiments) cannot be generalized from simple, concrete circuits. CertiQ may also fail to prove an optimization correct, e.g., because of complicated control code; in this case it falls back to translation validation, which adds extra cost and the possibility of failure at run-time. By contrast, every optimization in voqc has been proved correct. Finally, CertiQ does not directly represent the semantics of quantum programs, so it cannot be used as a tool for verifying general properties of a program’s semantics (as we do in Section 3 and Appendix B).

8 CONCLUSIONS AND FUTURE WORK

This paper has presented voqc, the first verified optimizer for quantum circuits implemented within a proof assistant. A key component of voqc is sqir, a simple, low-level quantum language deeply embedded in the the Coq proof assistant, which gives a semantics to quantum programs that is amenable to proof. Compiler passes are expressed as Coq functions which are proved to preserve the semantics of their input sqir programs. voqc’s optimizations are mostly based on local circuit equivalences, implemented by replacing one pattern of gates with another, or commuting a gate rightward until it can be cancelled. Others, like rotation merging, are more complex. These were inspired by, and in some cases generalize, optimizations in industrial compilers, but in voqc are proved correct. When applied to a benchmark suite of 28 circuit programs, we found voqc performed comparably to state-of-the-art compilers, reducing gate count on average by 18.7% compared to 10.4% for IBM’s Qiskit compiler, 10.9% for CQC’s t\ket compiler, and 26.4% for the cutting-edge research compiler by Nam et al. [2018]. Furthermore, voqc reduced $T$-gate count on average by 42.3% compared to 40.9% by Amy et al. [2013], 42.3% by Nam et al., and 43.8% by the PyZX compiler.

Moving forward, we plan to incorporate voqc into a full-featured verified compilation stack for quantum programs, following the vision of a recent Computing Community Consortium report [Martonos and Roetteler 2019]. We can implement validated parsers [Jourdan et al. 2012] for languages like OpenQASM and verify their translation to sqir (e.g., using metaQASM’s semantics [Amy 2019]); this work is already in progress [Singhal et al. 2020]. We can also add support for hardware-specific transformations that compile to a particular gate set. Indeed, most of the sophisticated code in Qiskit is devoted to efficiently mapping programs to IBM’s architecture, and IBM’s 2018 Developer Challenge centered around designing new circuit mapping algorithms [Staff 2018]. We leave it as future work to incorporate optimizations and mapping algorithms from
additional compilers into voqc. Our experience so far makes us optimistic about the prospects for doing so successfully.

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381–392.


When we first set out to build voqc, we thought to do it using Qwire [Paykin et al. 2017], another formally verified quantum programming language embedded in Coq. However, we were surprised to find that we had tremendous difficulty proving that even simple transformations were correct. This experience led to the development of sqir, and raised the question: Why does sqir seem to make proofs easier, and what do we lose by using it rather than Qwire?

The fundamental difference between sqir and Qwire is that sqir programs use concrete (numeric) indices into a global register to refer to qubits. As such, the semantics can naturally map qubits to rows and columns in the denoted matrix. In addition, qubit disjointness in a sqir program is obvious—$G_1 \ m$ operates on a different qubit than $G_2 \ n$ when $m \neq n$. Both elements are important for easily proving equivalences, e.g., that gates acting on disjoint qubits commute (a property that allows us to reason about gates acting on different parts of the circuit in isolation). In Qwire, variables are implemented using higher-order abstract syntax [Pfenning and Elliott 1988] and refer to abstract qubits. This approach eases programmability—larger circuits can be built by composing smaller ones, connecting inputs and outputs by normal variable binding, indifferent to the physical identity of a qubit. This approach is also used in the language Quipper [Green et al. 2013].

However, we find that this approach complicates formal proof. To denote the semantics of a program that uses abstract qubits requires deciding how abstract qubits will be represented concretely, as rows and columns in the denotation matrix. Reasoning about this translation can be laborious, especially for recursive circuits and those that allocate and deallocate qubits (entailing de Bruijn-style index shifting [Rand 2018]). Moreover, notions like disjointness are no longer obvious—$G_1 \ x$ and $G_2 \ y$ for variables $x \neq y$ may not be disjoint if $x$ and $y$ could be allocated to the same concrete qubit.

The trade off between the two approaches is most evident in how they support composition.

**Composition in Qwire.** Qwire circuits have the following form:

```ocaml
Inductive Circuit (w : WType) : Set :=
| output : Pat w → Circuit w
| gate : ∀ {w1 w2}, Gate w1 w2 → Pat w1 → (Pat w2 → Circuit w) → Circuit w
| lift : Pat Bit → (𝔹 → Circuit w) → Circuit w.
```

Patterns Pat type the variables in Qwire circuits; their type index $w$ corresponds to some collection of bits and qubits. The circuit output $p$ is simply a wire with the wire type associated with $p$. Note that this, like all Circuits, is an open term. The definition of gate takes in a Gate parameterized by an input type $w1$ and an output $w2$, an appropriately typed input pattern, and a continuation of the form $Pat w2 → Circuit w$, which describes how the gate’s output is used in the rest of the circuit. Finally, lift takes a single Bit (or classical wire) and a continuation that constructs a circuit based on that bit’s interpretation as a Boolean value.

This use of continuations makes composition easy to define:

```ocaml
Fixpoint compose {w1 w2} (c : Circuit w1) (f : Pat w1 → Circuit w2) : Circuit w2 :=
  match c with
  | output p ⇒ f p
  | gate g p c' ⇒ gate g p (fun p' ⇒ compose (c' p') f)
  | lift p c' ⇒ lift p (fun bs ⇒ compose (c' bs) f)
end.
```

In each case, the continuation is applied directly to the output of the first circuit.

While circuits correspond to open terms, closed terms are represented by boxed circuits:
Inductive Box w1 w2 : Set := box : (Pat w1 \to Circuit w2) \to Box w1 w2.

For convenience, box (fun w \to c) is written as simply box w \to c and let p \leftarrow c1 ; c2 is similarly defined as notation for compose c1 (fun p \to c2). One can unbox a boxed circuit Box w1 w2 simply by providing a valid Pat w1, obtaining a Circuit w2.

This representation allows for easy sequential and parallel composition of closed circuits. Running two circuits in sequence involves connecting the output of the first circuit to the input of the second circuit (where the types will guarantee compatibility) and running them in parallel gives the circuits disjoint inputs and outputs their results (leading to a tensor type).

Definition inSeq {w1 w2 w3} (c1 : Box w1 w2) (c2 : Box w2 w3) : Box w1 w3 :=
  box p1 \Rightarrow
  let p2 ← unbox c1 p1;
  unbox c2 p2.

Definition inPar {w1 w2 w1' w2'} (c1 : Box w1 w2) (c2 : Box w1' w2') :
  Box (w1 \otimes w1') (w2 \otimes w2') :=
  box (p1,p2) \Rightarrow
  let p1' ← unbox c1 p1;
  let p2' ← unbox c2 p2;
  (p1',p2').

Unfortunately, proving useful specifications for these functions is quite difficult. The denotation of a circuit is (in the unitary case) a square matrix of size $2^n$ for some $n$. To construct such a matrix we need to map all of a circuit’s (abstract) variables to 0 through $n-1$, ensuring that the mapping function has no gaps even when we initialize or discard qubits. Qwire maintains this invariant through compiling to a de Bruijn-style variable representation [de Bruijn 1972]. Reasoning about the denotation of circuits, then, involves reasoning about this compilation procedure. In the case of open circuits (the most basic circuit type), we must also reason about the contexts that type the available variables, which change upon every gate application.

Composition in sqir. Composing two sqir programs requires manually defining a mapping from the global registers of both programs to a new, combined global register. To do this, we provide two helper functions, which respectively renumber a unitary program’s concrete indexes according to a mapping $f$, and change the program’s global register size.

Fixpoint map_qubits {U dim} (f : \mathbb{N} \to \mathbb{N}) (c : ucom U dim) : ucom U dim :=
  match c with
  | c1; c2 ⇒ map_qubits f c1; map_qubits f c2
  | uapp1 u n ⇒ uapp1 u (f n)
  | uapp2 u m n ⇒ uapp2 u (f m) (f n)
  end.

Fixpoint cast {U dim} (c : ucom U dim) dim' : ucom U dim' :=
  match c with
  | c1; c2 ⇒ cast c1 dim' ; cast c2 dim'
  | uapp1 u n ⇒ uapp1 u n
  | uapp2 u m n ⇒ uapp2 u m n
  end.

With these, we can define parallel composition in sqir:
Definition \( \text{inPar} \{U \text{ dim1 dim2}\} \ (c1 : \text{ ucom} U \text{ dim1}) \ (c2 : \text{ ucom} U \text{ dim2}) := \)
\[
\begin{align*}
\text{cast} \ (\text{c1}) \ (\text{dim1} + \text{dim2}); \\
\text{cast} \ (\text{map_qubits} (\text{fun} \ q \Rightarrow \text{dim1} + q) \ c2) \ (\text{dim1} + \text{dim2}).
\end{align*}
\]

The correctness property for \( \text{inPar} \) says that the denotation of \( \text{inPar} \ c1 \ c2 \) can be constructed from the denotations of \( c1 \) and \( c2 \).

Lemma \( \text{inPar\_correct} \) : \( \forall \ c1 \ c2 \ d1 \ d2, \)
\[
\text{uc\_well\_typed} \ d1 \ c1 \rightarrow \[\text{inPar} \ c1 \ c2 \ d1 \ d2\] \ = \ [c1]_{d1} \otimes [c2]_{d2}.
\]

The \( \text{inPar} \) function is relatively simple, but more involved than the corresponding \( \text{Qwire} \) definition because it requires relabeling the qubits in program \( c2 \).

General composition in \( \text{sqir} \) (including sequential composition) requires even more involved relabeling functions that are less straightforward to describe. For example, consider the composition expressed in the following \( \text{Qwire} \) program:
\[
\text{box} \ (ps, q) \Rightarrow \\
\text{let} \ (x, y, z) \leftarrow \text{unbox} \ c1 \ ps; \\
\text{let} \ (q, z) \leftarrow \text{unbox} \ c2 \ (q, z); \\
(x, y, z, q).
\]

This program connects the last output of program \( c1 \) to the second input of program \( c2 \). This operation is natural in \( \text{Qwire} \), but describing this type of composition in \( \text{sqir} \) requires some effort. In particular, the programmer must determine the required size of the new global register (in this case, \( 4 \)) and explicitly provide a mapping from qubits in \( c1 \) and \( c2 \) to indices in the new register (for example, the first qubit in \( c2 \) might be mapped to the fourth qubit in the new global register).

When \( \text{sqir} \) programs are written directly, this puts extra burden on the programmer. When \( \text{sqir} \) is used as an intermediate representation, however, these mapping functions should be produced automatically by the compiler. The issue remains, though, that any proofs we write about the result of composing \( c1 \) and \( c2 \) will need to reason about the mapping function used (whether produced manually or automatically).

As an informal comparison of the impact of \( \text{Qwire} \)’s and \( \text{sqir} \)’s representations on proof, we note that while proving the correctness of the \( \text{inPar} \) function in \( \text{sqir} \) took a matter of hours, there is no correctness proof for the corresponding function in \( \text{Qwire} \), despite many months of trying. Of course, this comparison is not entirely fair: \( \text{Qwire} \)’s \( \text{inPar} \) is more powerful than \( \text{sqir} \)’s equivalent.

\( \text{sqir} \)’s \( \text{inPar} \) function does not require every qubit within the global register to be used – any gaps will be filled by identity matrices. Also, \( \text{sqir} \) does not allow introducing or discarding qubits, which we suspect will make ancilla management difficult to reason about.

Quantum Data Structures. \( \text{sqir} \) also lacks some other useful features present in higher-level languages. For example, in QIO [Altenkirch and Green 2010] and Quipper [Green et al. 2013] one can construct circuits that compute on quantum data structures, like lists and trees of qubits. In \( \text{Qwire} \), this concept is refined to use more precise dependent types to characterize the structures; e.g., the type for the \( \text{GHZ} \) program indicates it takes a list of \( n \) qubits to a list of \( n \) qubits. More interesting dependently-typed programs, like the quantum Fourier transform, use the parameter \( n \) as an argument to rotation gates within the program.

Regrettably, these structures can make reasoning about programs difficult. For instance, as shown in Figure 12, the \( \text{GHZ} \) program written in \( \text{Qwire} \) emits a list of qubits while the \( \text{fredkin\_seq} \) circuit takes in a tree of qubits. Connecting the qubits from a \( \text{GHZ} \) to a \( \text{fredkin\_seq} \) circuit with the same arity requires an intermediate gadget. And if we want to verify a property of this composition, we need to prove that this gadget is an identity. In \( \text{sqir} \), which has neither quantum data structures nor typed circuits, this issue does not present itself.
Fig. 12. The patterns for the output of ghz (left) and the input to fredkin_seq (right) over four qubits. ghz produces a list of qubits ((q1,q2),q3,q4) whereas fredkin_seq expects a tree ((q1,q2),(q3,q4)). In composition, the mismatched patterns require an extra gadget to transform the former into the latter.

**Dynamic Lifting.** sqir also does not support dynamic lifting, which refers to a language feature that permits measuring a qubit and using the result as a Boolean value in the host language to compute the remainder of a circuit [Green et al. 2013]. Dynamic lifting is used extensively in Quipper and QWIRE. Unfortunately, its presence complicates the denotational semantics, as the semantics of any Quipper or QWIRE program depends on the semantics of Coq or Haskell, respectively. In giving a denotational semantics to QWIRE, Paykin et al. [2017] assume an operational semantics for an arbitrary host language, and give a denotation for a lifted circuit only when both of its branches reduce to valid QWIRE circuits.

Although sqir does not support dynamic lifting, its meas construct is a simpler alternative. Since the outcome of the measurement is not used to compute a new circuit, sqir does not need a classical host language to do computation: It is an entirely self-contained, deeply embedded language. As a result, we can reason about sqir circuits in isolation, and also easily reason about families of sqir circuits described in Coq.

**Other Differences.** Another important difference between QWIRE and sqir is that QWIRE circuits cannot be easily decomposed into subcircuits because output variables are bound in different places throughout the circuit. By contrast, a sqir program is an arbitrary nesting of smaller programs. This means that for sqir, the program $c1;((c2;(c3;c4));c5)$ is equivalent to $c1;c2;c3;c4;c5$ under all semantics, whereas every QWIRE circuit (only) associates to the right. This allows us to arbitrarily flatten sqir programs into a convenient list representation, as is done in voqc (Section 4), and makes it easy to rewrite using sqir identities.

Also, unlike most quantum languages and as already discussed in the main body of the paper, sqir features a distinct core language of unitary operators; the full language adds measurement to this core. The semantics of a unitary program is expressed directly as a matrix, which means that proofs of correctness of unitary optimizations (the bulk of voqc) involve reasoning directly about matrices. Doing so is far simpler than reasoning about functions over density matrices, as is required for the full sqir language or any program in QWIRE.

**Concluding thoughts.** Upon reflection, we can see that the differences between QWIRE and sqir ultimately stem from their design goals. QWIRE was developed as a general-purpose programming language for quantum computers [Paykin et al. 2017], with ease of programmability as a key concern; only later was it extended as a tool for verification [Rand et al. 2018b, 2017]. By contrast, sqir was designed from the start with verification in mind, with by-hand programmability a secondary consideration; we expected sqir would be compiled from another language such as Q# [Svore et al. 2018], Quipper [Green et al. 2013] or even QWIRE itself. That said, for near-term quantum programs sqir’s lower-level abstractions have not proved difficult to use, even for source programming (as shown in the Appendix B). As programs scale up, finding the right way to extend sqir (or a language like it) with higher level abstractions without overly complicating verification will be an important goal.
B SQIR FOR PROGRAM VERIFICATION

sqir’s simple structure and semantics allow us to easily verify general properties of quantum programs. In this section we discuss correctness properties of three quantum programs written in sqir: superdense coding, the Deutsch-Jozsa algorithm, and quantum phase estimation. The Deutsch-Jozsa algorithm is verified for any Boolean oracle with any number of qubits and quantum phase estimation is verified for any input program and eigenvector with any number of qubits, demonstrating that sqir supports parameterized proofs. We also present an alternative proof for quantum teleportation using a non-deterministic semantics; the original is found in Section 3, along with our proof of GHZ state preparation.

B.1 Superdense Coding

Superdense coding is a protocol that allows a sender to transmit two classical bits (i.e., Booleans), $b_1$ and $b_2$, to a receiver using a single quantum bit. Doing so relies on the sender and receiver sharing a Bell pair, as shown in Figure 13. The sender conditionally applies X and Z to her qubit, contingent on the values of her bits, and then transmits it to the receiver, who applies a Bell measurement (reversed entangling operation followed by measure) to recover the original bits. The sqir program corresponding to the unitary part of this circuit is produced by the Coq function superdense, shown in Figure 14, which is parametrized by the classical input bits $b_1$ and $b_2$.

We can prove that the result of evaluating the program superdense $b_1 b_2$ on an input state consisting of two qubits initialized to zero is the state $\bigotimes^{b_1, b_2} \bigotimes^{\psi}$. Lemma superdense_correct : $\forall b_1 b_2$, $\parallel$ superdense $b_1 b_2 \parallel_d = |b_1, b_2\rangle$.

The proof simply destructs $b_1$ and $b_2$ and applies our matrix simplification tactics.

B.2 Teleport with Nondeterministic Semantics

The proof of correctness of quantum teleportation using the density matrix semantics (Section 3.2) is simple, but not particularly useful for understanding why the protocol is correct. A more illuminating proof can be carried out using an alternative nondeterministic semantics in which evaluation is expressed as a relation. Given a state $\langle\psi\rangle$, unitary program $u$ will (deterministically) evaluate to $\parallel u_d \times \langle\psi\rangle \parallel$. However, meas $q p_1 p_2$ may evaluate to either $p_1$ applied to $|1angle\langle 1| \times \langle\psi\rangle$ or $p_2$ applied to $|0\rangle\langle 0| \times \langle\psi\rangle$. At every point in the program, the nondeterministic semantics represents the state using a vector $\langle\psi\rangle$ rather than a density matrix $\rho$. 
Fixpoint npar n (u : N → ucom base n) :=
match n with
| 0 ⇒ SKIP
| S n' ⇒ npar n' u ; u n'
end.

Definition deutsch_jozsa n (u : ucom base n) :=
X (n-1) ; npar n H ; u ; npar n H.

Fig. 15. The Deutsch-Jozsa algorithm in sqir and as a circuit. The Coq function npar constructs a sqir program that applies the same operation to every qubit.

However, because we do not track all possible measurement outcomes, the nondeterministic semantics is only useful for proving properties for which all outcomes lead to the same result. The correctness of the teleport protocol is an example of such a property, because we obtain the original qubit regardless of the measurement outcome. Similarly, the nondeterministic semantics is only useful for verifying transformations of programs involving deterministic outcomes. In practice, we use the density matrix semantics to verify our optimizations.

Using the non-deterministic semantics, the proof of teleport is more involved, but also more illustrative of the inner workings of the algorithm. Under the non-deterministic semantics, we aim to prove the following:

Lemma teleport_correct :
∀ (ψ : Vector (2^1)) (ψ′ : Vector (2^3)),
WF_Matrix ψ → teleport / (ψ ⊗ |0,0⟩) ↓ ψ′ → ψ′ ∝ |0,0⟩ ⊗ ψ.

This says that on input |ψ⟩⊗|0,0⟩, teleport will produce a state that is proportional (∝) to |0,0⟩⊗|ψ⟩. Note that this statement is quantified over every outcome ψ′ and hence all possible paths to ψ′. If instead we simply claimed that teleport / (ψ ⊗ |0,0⟩) ↓ 1/2 · (|0,0⟩ ⊗ ψ), where the 1/2 factor reflects the probability of each measurement outcome ((1/2)^3 = 1/4), we would only be stating that some such path exists.

The first half of the circuit is unitary, so the proof simply computes the effect of applying a H gate, two CNOT gates and another H gate to the input vector state. The two measurement steps then leave us with four different cases to consider. In each of the four cases, we can use the outcomes of measurement to correct the final qubit, putting it into the state |ψ⟩. Finally, resetting the already-measured qubits is deterministic and leaves us with the desired state.

B.3 Deutsch-Jozsa Algorithm

In the Deutsch-Jozsa problem [Deutsch and Jozsa 1992], the goal is to determine whether a Boolean function f : {0,1}^n → {0,1} is constant (always returns the same value) or balanced (returns 0 and 1 equally often), given that one is the case. The function f is encoded in an oracle U : |x, y⟩ ↦ |x, y ⊕ f(x)⟩, which is a linear operator over a 2^{n+1} dimensional Hilbert space. In Coq, we express the requirement that program u encodes the function f as follows.

Definition boolean_oracle {n} (u : ucom (n + 1)) f :=
∀ x, y, [u]_n+1 × |x⟩ ⊗ |y⟩ = |x⟩ ⊗ |y ⊕ (f x)⟩.

To express that a function is constant or balanced, we can define a function count f n that counts all inputs on which function f (with domain size 2^n) evaluates to true. Then we have:

Definition balanced f n := n > 0 ∧ count f n = 2^{n-1}.
Definition constant f n := count f n = 0 ∨ count f n = 2^n.
As shown in Figure 15, the Deutsch-Jozsa algorithm begins with an all $|0\rangle$ state and prepares the input state $|+\rangle^{\otimes n} \otimes |-\rangle$ by applying an $X$ gate on the last qubit followed by an $H$ gate on every qubit. Next the oracle $U$ is queried, and a $H$ gate is again applied to every qubit in the program. Finally, all qubits are measured in the standard basis. If measuring all the qubits but the last yields an all-zero string (the last qubit is guaranteed to be in the $|1\rangle$ state) then the algorithm outputs “accept,” indicating that the function is constant. Otherwise the algorithm outputs “reject.”

The probability of measuring $|0\rangle$ in the first $n$ qubits and $|1\rangle$ in the last qubit is given by

$$Pr_{\text{accept}} = |\langle 0 |^{\otimes n} \otimes \langle 1 | \rangle \times (\text{deutsch\_jozsa} n U \times |0\rangle^{\otimes(n+1)})|^2.$$ 

We define an accept predicate that states that $Pr_{\text{accept}} = 1$ and a reject predicate that states that $Pr_{\text{accept}} = 0$. Our correctness property is then stated as follows.

**Lemma** deutsch\_jozsa\_correct : 
\[
\forall (n : \mathbb{N}) \ (f : \mathbb{N} \rightarrow \mathbb{B}) \ (u : \text{base\_ucom} (n + 1)), \\
n > 0 \rightarrow \text{boolean\_oracle} \ u \ f \rightarrow \\
(\text{constant} \ f \ n \rightarrow \text{accept} \ u) \land (\text{balanced} \ f \ n \rightarrow \text{reject} \ u).
\]

The key lemma in our proof states that $Pr_{\text{accept}}$ depends on the number of inputs on which $f$ evaluates to 1, i.e., count $f$ $n$. In particular, $Pr_{\text{accept}} = |1 - \frac{2\times\text{count} \ f \ n}{2^{n}}|^2$. We prove this property using matrix simplification and induction on $n$. Correctness of the Deutsch-Jozsa algorithm follows directly from this lemma. For a constant function, count $f$ $n = 0$ or count $f$ $n = 2^n$ so $Pr_{\text{accept}} = 1$. For a balanced function, count $f$ $n = 2^{n-1}$ so $Pr_{\text{accept}} = 0$.

Unlike the GHZ state preparation, quantum teleportation, and superdense coding examples presented so far, the deutsch\_jozsa program is parameterized by both the size of the global register $n$ and a sqir program $u$. Many quantum algorithms are described as such families of circuits, so parameterized programs are an important target for verification tools [Charetton et al. 2020]. Proofs about parameterized programs are only possible in tools that manipulate the semantics matrix symbolically.

### B.4 Quantum Phase Estimation

Given a unitary matrix $U$ and eigenvector $|\psi\rangle$ such that $U |\psi\rangle = e^{2\pi i \theta} |\psi\rangle$, the goal of quantum phase estimation (QPE) is to find a $k$-bit representation of $\theta$. In the case where $\theta$ can be exactly represented using $k$ bits (i.e. $\theta = z/2^k$ for some $z \in \mathbb{Z}$), QPE recovers $\theta$ completely. Otherwise, the algorithm finds a good $k$-bit approximation with high probability. QPE is often used as a subroutine in quantum algorithms, most famously Shor’s factoring algorithm [Shor 1994]. For more details on phase estimation see Nielsen and Chuang [2000, Chapter 5].

The circuit for QPE is shown in Figure 16. First, a layer of Hadamard gates prepares a uniform superposition. Next, a sequence of controlled $U$ operations encodes information about $\theta$ in the phase. Finally, the inverse quantum Fourier transform (QFT) is used to recover the information about $\theta$ stored in the phase. The full sqir definition of QPE is shown in Figure 17.

The statement of correctness for the case where $\theta$ can be described exactly using $k$ bits ($\theta = z/2^k$) is as follows.

**Lemma** QPE\_semantics\_simplified : $\forall \ k \ n \ (u : \text{uc\_well\_typed} \ n) \ z \ (\psi : \text{Vector} \ 2^n)$,
\[
n > 0 \rightarrow k > 1 \rightarrow \text{uc\_well\_typed} \ u \rightarrow \text{WF\_Matrix} \ \psi \rightarrow \\
\text{let} \ \theta := \frac{z}{2^k} \ \text{in} \\
\|u\|_n \times \psi = \frac{e^{2\pi i \theta}}{*} \psi \rightarrow \\
\|\text{QPE} \ k \ n \ u\|_{k+n} \times (|0\rangle^{\otimes k} \otimes \psi) = |z\rangle \otimes \psi.
\]
\[ QPE_{k,n} = \begin{array}{c}
|0\rangle \quad H \\
\vdots \\
|0\rangle \quad H \\
|\psi\rangle \quad U^n \quad U^2 \quad \cdots \quad U^{2^k+1} \\
\end{array} \begin{array}{c}
QFT^{-1}_k \\
\vdots \\
H \\
R_{k-2} \\
H \\
R_{k-1} \\
R_k \\
R_{k-1} \\
H \\
R_2 \\
H \\
\end{array} \]

The first four conditions ensure well-formedness of the inputs. The fifth condition enforces that input \( \psi \) is an eigenvector of \( c \). The conclusion says that running the QPE program computes the value \( z \), as desired. Note that like the deutsch_jozsa_correct, this property is parameterized by a \( s_qir \) program.

In the case where \( \theta \) can not be exactly described using \( k \) bits, we instead prove that QPE recovers the best \( k \)-bit approximation with high probability (in particular, with probability \( \geq \frac{4}{\pi^2} \)).

**Lemma QPE_semantics_full:**

\[
\forall \; k \; n \; (u : \text{uc base n}) \; z \; (\psi : \text{Vector} \; 2^n) \; (\delta : \mathbb{R}), \\
\quad n > 0 \rightarrow k > 1 \rightarrow \text{uc_well_typed} \; u \rightarrow \text{Pure_State_Vector} \; \psi \rightarrow \\
\quad -1 \; / \; 2^{k+1} < \delta < 1 \; / \; 2^{k+1} \rightarrow \delta \neq 0 \rightarrow \\
\quad \text{let } \theta := z / 2^k + \delta \text{ in} \\
\quad \|u\|_n \times \psi = e^{2\pi i \theta} \star \psi \rightarrow \\
\quad \text{probability_of_outcome} \; (\|QPE \; k \; n \; u\|_{k+n} \times (|0\rangle^{\otimes k} \otimes \psi)) \; (|z\rangle \otimes \psi) \; \geq \; \frac{4}{\pi^2}.
\]

As above, \( \text{probability_of_outcome} \; |\psi\rangle \; |\phi\rangle \) is the probability of measuring outcome \( |\phi\rangle \) given input \( |\psi\rangle \). Pure_State_Vector is more restrictive form of WF_Matrix that requires a vector to have norm 1.  

\( s_qir \)’s proofs of correctness for QPE rely on a variety of proofs about its components. For example, we prove that \( \text{ncar} \; k \; H \) computes a uniform superposition \( \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} |k\rangle \) when given input \( |0\rangle^{\otimes k} \) and \( \text{QFT} \; n \) maps the state \( |x\rangle \) to \( \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} e^{2\pi i k/x} |k\rangle \). Our proofs typically proceed by induction on the dimension variable and make heavy use of our matrix simplification tactics. In addition, the proof for \( \text{QPE_semantics_full} \) relies on several lemmas bounding the behavior of trigonometric functions. These are proved correct using the Coq Interval package [Melquiond 2020] and its associated tactic.¹

**C ADDITIONAL BENCHMARK RESULTS**

\( \text{voqc} \) is able to run on all 99 benchmark programs considered by Nam et al. [2018]. In this section, we compare \( \text{voqc} \)’s performance against Nam et al.’s on these benchmarks, confirming our conclusion from Section 6: \( \text{voqc} \) is a faithful implementation of a subset of the optimizations present in Nam et al., which according to Section 6’s results is the most effective of several compilers we

¹Thanks to Laurent Théry for the pointer to Interval and providing proofs of our sin_sublinear and sin_PIx_ge_2x lemmas.
A Verified Optimizer for Quantum Circuits

(* Controlled rotation cascade on n qubits. *)

Fixpoint controlled_rotations n : ucom base n :=
  match n with
  | 0 | 1 => SKIP
  | S n' => controlled_rotations n'; control n' (Rz (2π / 2^n) 0)
  end.

(* Quantum Fourier transform on n qubits. 'map_qubits' applies a function to all qubit references. In this case it increments all qubit references by 1. *)

Fixpoint QFT n : ucom base n :=
  match n with
  | 0 => SKIP
  | 1 => H 0
  | S n' => H 0 ; controlled_rotations n ; map_qubits (fun q => q + 1) (QFT n')
  end.

(* QFT outputs qubits in the wrong order, so the qubits need to be reversed before further processing. This can be handled by the classical control hardware or on the quantum machine with SWAPs, as done here. *)

Fixpoint reverse_qubits' dim n : ucom base dim :=
  match n with
  | 0 => SKIP
  | S n' => reverse_qubits' dim n' ; SWAP n' (dim - n' - 1)
  end.

Definition reverse_qubits n := reverse_qubits' n (n/2).
Definition QFT_w_reverse n := QFT n ; reverse_qubits n.

(* Controlled powers of u *)

Fixpoint controlled_powers' {n} (u : ucom base n) k kmax : ucom base (kmax + n) :=
  match k with
  | 0 => SKIP
  | S k' => controlled_powers' u k' kmax ;
    niter 2^k' (control (kmax - k' - 1) u)
  end.

Definition controlled_powers {n} (u : ucom base n) k := controlled_powers' u k k.

(* QPE circuit for program u. k = number of bits in resulting estimate n = number of qubits in input state *)

Definition QPE k n (u : ucom base n) : ucom base (k + n) :=
  npar k H ;
  controlled_powers (map_qubits (fun q => k + q) u) k;
  invert (QFT_w_reverse k).

Fig. 17. sqir definition of QPE. Some type annotations and calls to cast (a no-op that changes the program dimension) have been removed for clarity. control, map_qubits, niter, and invert are Coq functions that transform sqir programs; we have proved that they have the expected behavior (e.g. ⟦invert u⟧_n = ⟦u⟧_n^†) for any input program.
experimented with on this benchmark set. Our versions of the benchmarks are available online. All results were obtained using a laptop with a 2.9 GHz Intel Core i5 processor and 16 GB of 1867 MHz DDR3 memory, running macOS Mojave. Nam et al.’s results are from a similar machine with 8 GB RAM running OS X El Capitan. Their implementation is written in Fortran. In our tables we report both the time required for \texttt{voqc} to parse the input file and to perform optimization.

Overall, the results are consistent with those presented in Section 6. In cases where Toffoli decomposition and heavy optimization are not used (the QFT, QFT-based adder, and product formula circuits), \texttt{voqc}’s results are identical to Nam et al.’s. In the other cases, \texttt{voqc} is slightly less effective than Nam et al. for the reasons discussed in Section 6. This is strong evidence that our implementation is faithful to Nam et al.’s (along with being proved correct!). In the worst case, \texttt{voqc}’s run time is four orders of magnitude worse than Nam et al.’s. However, \texttt{voqc}’s run time is often less than a second. We view this performance as acceptable, given that benchmarks with more than 1000 two-qubit gates (the only programs for which \texttt{voqc} optimization takes longer than one second) are well out of reach of current quantum hardware [Preskill 2018]. We are optimistic that \texttt{voqc}’s performance can be improved through more careful engineering.

Nam et al.’s benchmarks are divided into three categories; we describe each category below.

\textbf{Arithmetic and Toffoli.} These benchmarks are a superset of the arithmetic and Toffoli circuits discussed in Section 6. They range from 45 to 346,533 gates and 5 to 489 qubits. Results on all 32 benchmarks are given in Table 4. As discussed in Section 6, \texttt{voqc}’s performance does not match Nam et al.’s because we have not yet implemented all of their transformations (in particular, we are missing “Toffoli decomposition” and “Floating Rz gates”).

\textbf{QFT and Adders.} These benchmarks consist of components of Shor’s integer factoring algorithm, in particular the quantum Fourier transform (QFT) and integer adders. Two types of adders are considered: an in-place modulo 2q adder implemented in the Quipper library and an in-place adder based on the QFT. These benchmarks range from 148 to 381,806 gates and 8 to 4096 qubits. Results on all 27 benchmarks are given in Table 5, Table 6, and Table 7. The Quipper adder programs use similar gates to the arithmetic and Toffoli circuits, so the results are similar—\texttt{voqc} is close to Nam et al., but under-performs due to our simplified Toffoli decomposition. The QFT circuits use rotations parameterized by $\pi/2^n$ for varying $n \in \mathbb{N}$ (and no Toffoli gates) so \texttt{voqc}’s results are identical to Nam et al.’s. For consistency with Nam et al., on the QFT and QFT-based adder circuits we run a simplified version of our optimizer that does not include rotation merging.

\textbf{Product Formula.} These benchmarks implement product formula algorithms for simulating Hamiltonian dynamics. The benchmarks range from 260 to 127,500 gates and 10 to 100 qubits; they use rotations parameterized by floating point numbers, which we convert to OCaml rationals at parse time. The product formula circuits are intended to be repeated for a fixed number of iterations, and our resource estimates account for this. \texttt{voqc} applies Nam et al.’s “LCR” optimization routine to optimize programs across loop iterations.

On all 40 product formula benchmarks, our results are the same as those reported by Nam et al. [2018, Table 3]. $H$ gate reductions range from 62.5% to 75%. Reductions in Clifford $z$-axis rotations (i.e. rotations by multiples of $\pi/2$) range from 75% to 87.5% while reductions in non-Clifford $z$-axis rotations range from 0% to 28.6%. \texttt{CNOT} gate reductions range from 0% to 33%. Runtimes range from 0.02s for parsing and optimizing to 635.89s for parsing and 398.47s for optimizing. By comparison, Nam et al.’s runtimes range from 0.004s to 0.137s.

\footnote{https://github.com/inQWIRE/SQIR/tree/master/VOQC/benchmarks}
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<th>Nam (L) t(s)</th>
<th>Nam (H) Total</th>
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| Avg. Red.        | 24.6%       | 26.4%         | 19.2%        |

Table 4. Total gate count reduction on the “Arithmetic and Toffoli” circuits. Includes all programs listed in Table 2 and Table 3 as well as four gf programs omitted from that table. The reported voqc time only includes optimization time. voqc’s parse time was less than 0.001s for all benchmarks except the larger gf programs; the largest, gf2^163_mult, required 463s (7.7min) to parse. Nam (H) results were not available for the large benchmarks.
Table 5. Total gate count reduction on Quipper adder circuits. \( \text{voqc}'s \) \( H \) and \( T \) counts are identical to Nam (L) and (H), but the total \( R_z \) and \( CNOT \) counts are higher due to Nam et al.'s specialized Toffoli decomposition. The difference between Nam (L) and Nam (H) is entirely due to \( CNOT \) count. Our initial gate counts are higher than those reported by Nam et al. because we do not have special handling for +/- control Toffoli gates; we simply consider the standard Toffoli gate conjugated by additional \( X \) gates.

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<tr>
<th>( n )</th>
<th>Original ( \text{Total} )</th>
<th>Nam (L) ( \text{Total} ) ( t(s) )</th>
<th>Nam (H) ( \text{Total} ) ( t(s) )</th>
<th>( \text{voqc} ) ( \text{Total} )</th>
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</table>

| Avg. Red. | 63.7% | 71.6% | 45.7% |

Table 6. Results on QFT circuits. Exact timings and gate counts are not available for Nam (L) or Nam (H), but our results are consistent with those reported in Nam et al. [2018, Figure 1].

<table>
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<th>Original ( \text{CNOT} ) ( R_z ) ( H )</th>
<th>( \text{voqc} ) ( \text{CNOT} ) ( R_z ) ( H )</th>
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| Avg. Red. | 0% | 59.3% | 0% |

Table 7. Results on QFT-based adder circuits. Final gate counts are identical for \( \text{voqc} \) and Nam (L).